

Percolation noise at the metal–insulator transition of nanostructured VO₂ films

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Measurements: (2008 unsuccessful) successful: 2012 – 2014

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Laszlo

Shuyi



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Claes



Resistance noise at the metal–insulator transition in thermochromic VO₂ films

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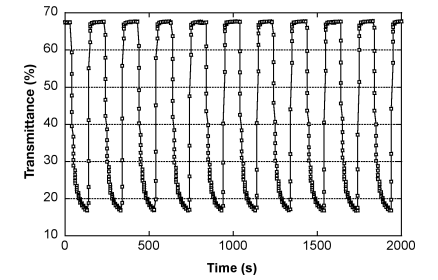
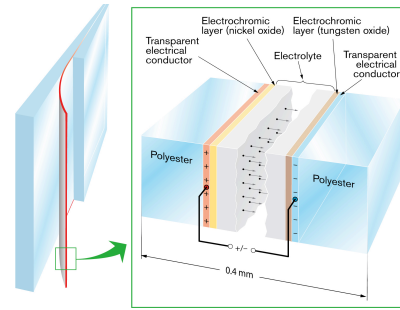
Thermochromic VO₂ films were prepared by reactive DC magnetron sputtering onto heated sapphire substrates and were used to make 100-nm-thick samples that were 10 μm wide and 100 μm long. The resistance of these samples changed by a factor ~2000 in the 50 < T_s < 70 °C range of temperature T_s around the “critical” temperature T_c between a low-temperature semiconducting phase and a high-temperature metallic-like phase of VO₂. Power density spectra S(f) were extracted for resistance noise around T_c and demonstrated unambiguous 1/f behavior. Data on S(10 Hz)/R_s² scaled as R_s^x, where R_s is sample resistance; the noise exponent x was –2.6 for T_s < T_c and +2.6 for T_s > T_c. These exponents can be reconciled with the Pennetta–Trefán–Reggiani theory [C. Pennetta, G. Trefán, and L. Reggiani, Phys. Rev. Lett. 85, 5238 (2000)] for lattice percolation with switching disorder ensuing from random defect generation and healing in steady state. Our work hence highlights the dynamic features of the percolating semiconducting and metallic-like regions around T_c in thermochromic VO₂ films.

VO₂: applications thermochromic glazing

Energy-Saving Applications. Smart-windows:
Voltage or temperature controlled transparency

Transparent

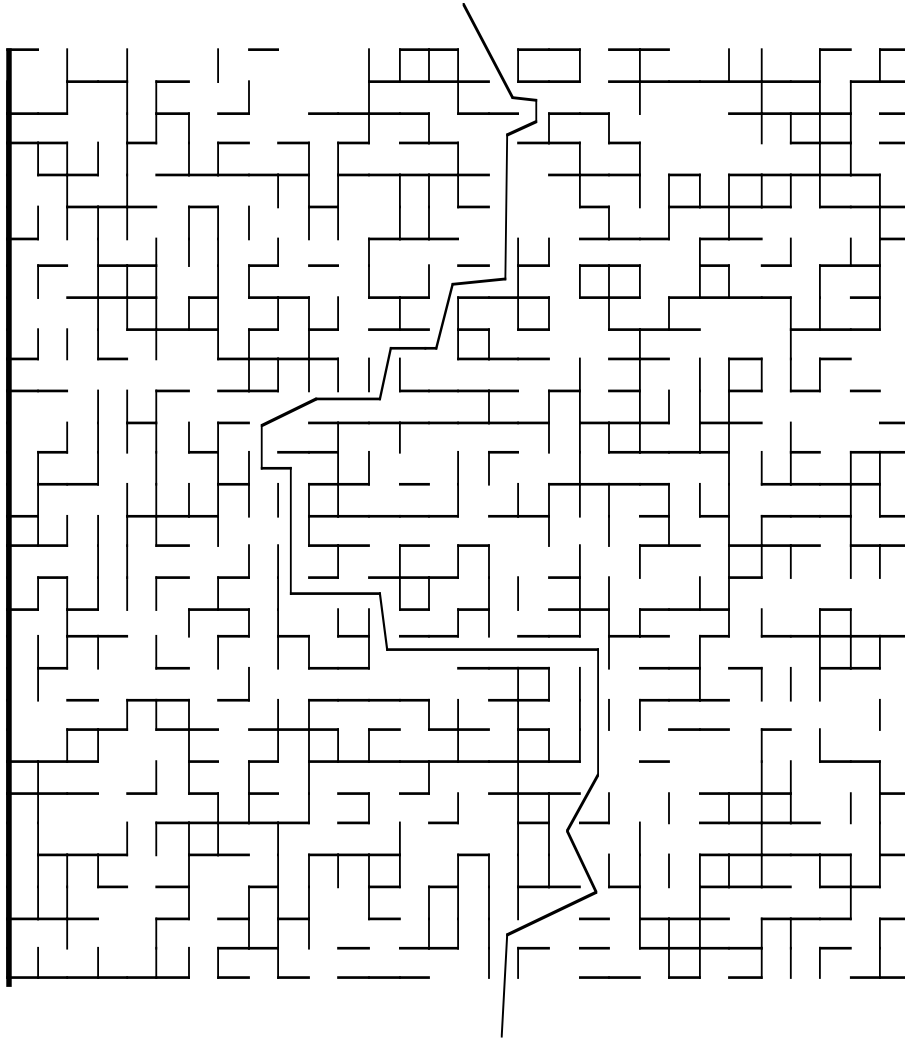
Darkened



VO₂:

- Thermochemical [F. J. Morin, Phys. Rev. Lett. 3, 34 (1959)];
- Single crystal: first-order metal–insulator transition (MIT) at $T_c \approx 68$ °C; switching between a low-temperature (monoclinic) *semiconducting* state and a high-temperature (rutile) *metallic-like* state.
- Thin films: the MIT is gradual with *metallic-like* regions growing in extent as the sample temperature T_s approaches T_c from below and with *semiconducting* regions disappearing as T_s becomes increasingly larger than T_c .

Lattice percolation



$$R(p) \propto (p - p_c)^{-t}$$

p : filling factor (0 – 1)

p_c : percolation threshold

t : resistivity exponent (1 - 2)

depends on dimension only (in regular lattices)

because percolation length is scaling in a similar fashion

$$\frac{S_R(f)}{R^2} \propto R^x$$

resistivity and noise exponent depends on dimension only (in regular lattices) and their absolute value is the same in 2D for conductor-insulator/superconductor transition (due to duality in 2D)

figure from: Z. Gingl, et al, Semicond. Sci. & Technol. 11 (1996) 1770.

Lattice percolation is confirmed in VO₂ films by J. Rozen, et al, Appl. Phys. Lett. 88, 081902 (2006).

p : determined from transparance measurements

$$R(p) \propto (p - p_c)^{-t}$$

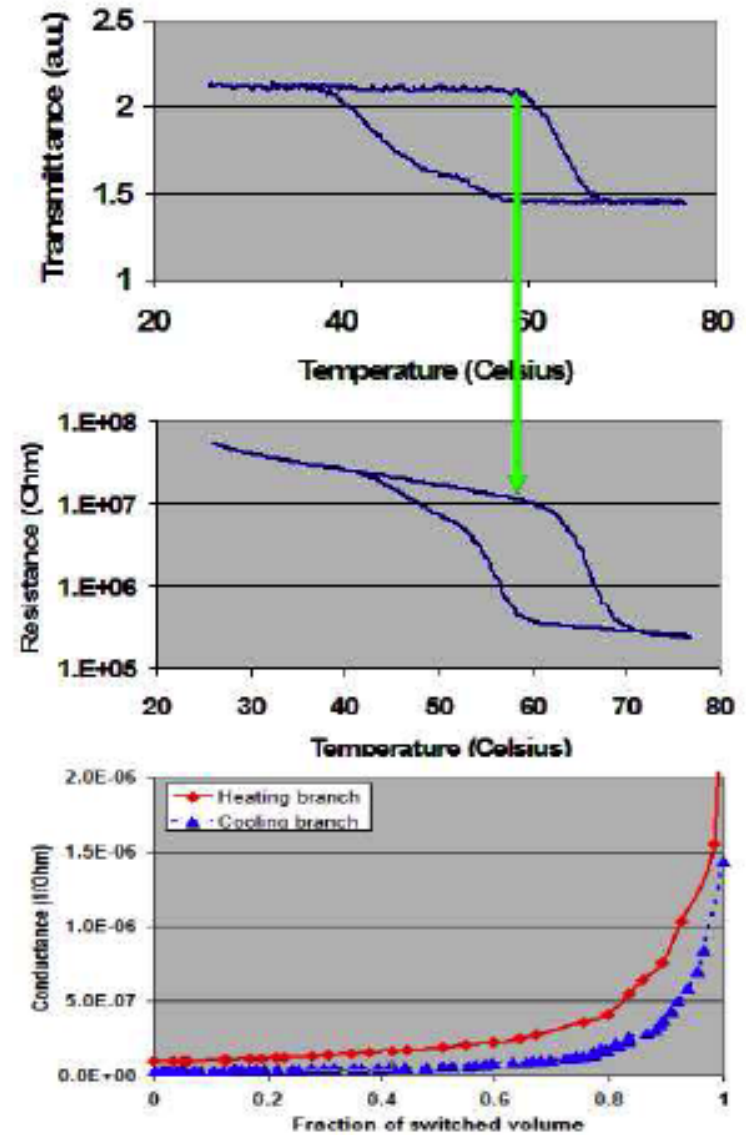
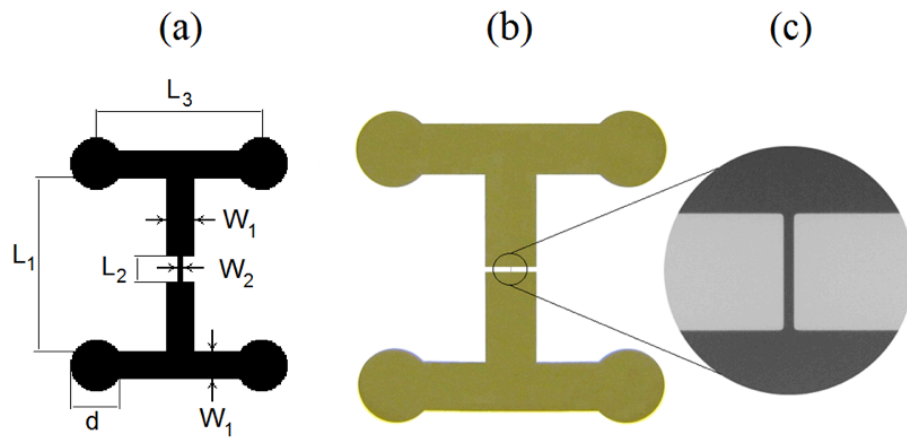


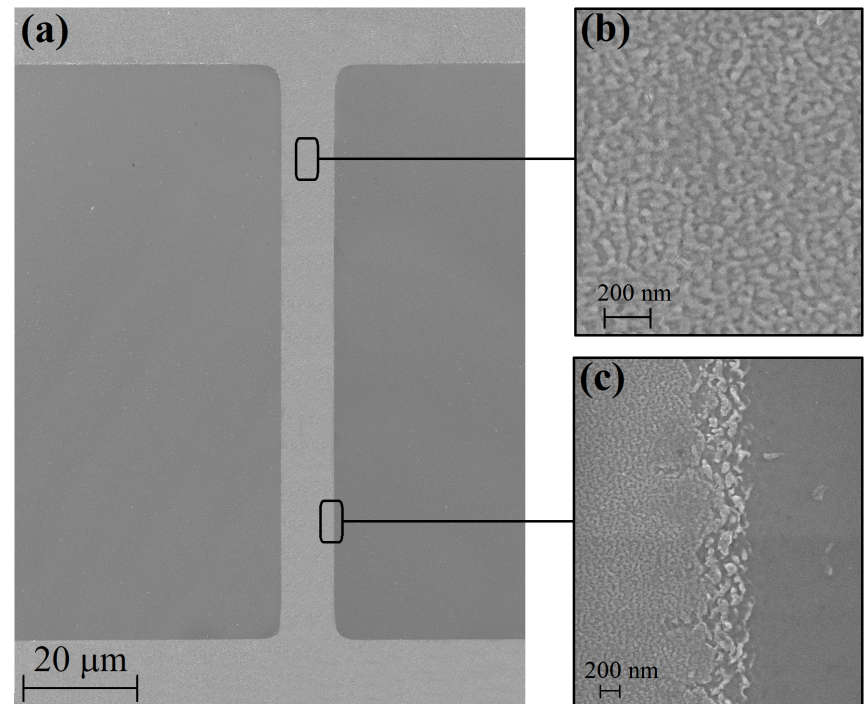
FIG. 3 Transmittance (a) and resistance (b) curves measured simultaneously as a function of temperature. (c) Dependence of the conductance on the volume fraction of the metallic phase deduced from the experimental data.

Noise experiments.

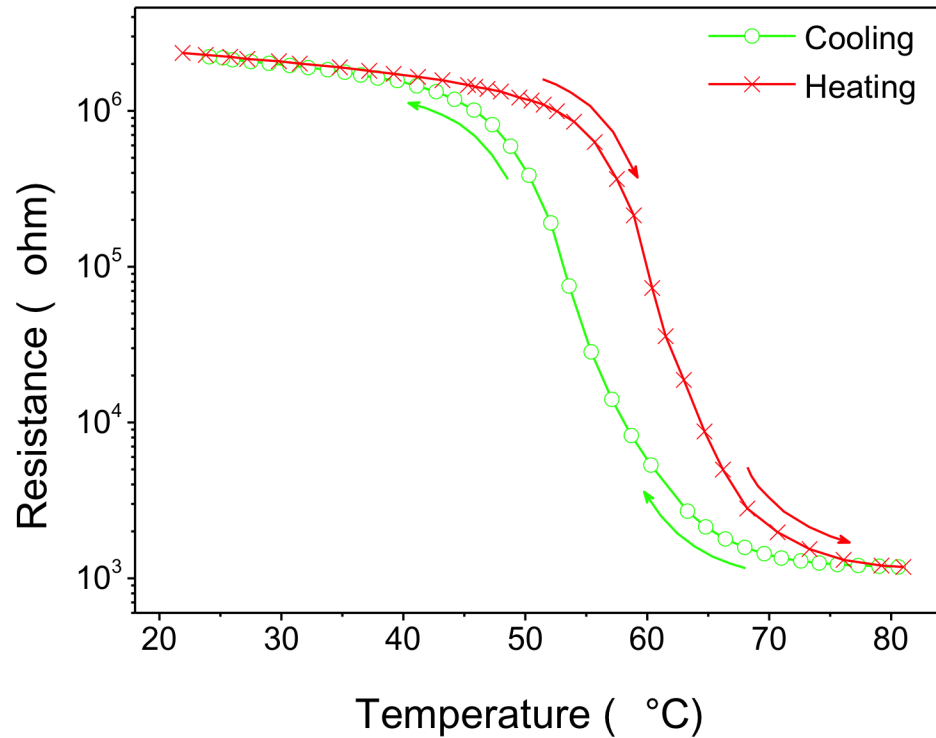
Schematic illustration of a **100-nm**-thick VO₂ sample with a micro-bridge in the middle and with contact pads for four-point electrical measurements. The dimensions are $L_1 = 5$ mm, $L_2 = \mathbf{100\ \mu\text{m}}$, $L_3 = 6$ mm, $d = 1.5$ mm, $W_1 = 1$ mm and $W_2 = \mathbf{10\ \mu\text{m}}$. Panels (b) and (c) are photos of the same structure and of the VO₂ micro-bridge in the encircled region, respectively.



SEM micrographs of the VO₂ micro-bridge



Resistance measurements



Extraordinary temperature sensitivity. For successful conductance noise measurements, it requires *ultra-low noise temperature control*, which is not available on the market.

We had similar problems in the 1990's with measuring high- T_c superconductor noise.

New Noise Exponents in Random Conductor-Superconductor and Conductor-Insulator Mixtures

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(Received 29 April 1993)

Time dependent fluctuations of the fraction of normal-conducting part in random resistor-superconductor (RS) and resistor-insulator (RI) networks lead to a novel effect close to the percolation threshold. The normalized noise scales as a function of the resistance with a characteristic exponent λ . The value of λ is different from the value found in classical percolation models but can be related to the resistivity exponent $s(t)$ of the RS (RI) transition by a simple scaling relation: $\lambda = 2/s$ ($2/t$). Results of recent experiments on high- T_c superconducting thin films are interpreted in terms of this new effect and a crossover from three to two dimensional percolation behavior is found.

September 19, 1989



Peter

Laszlo

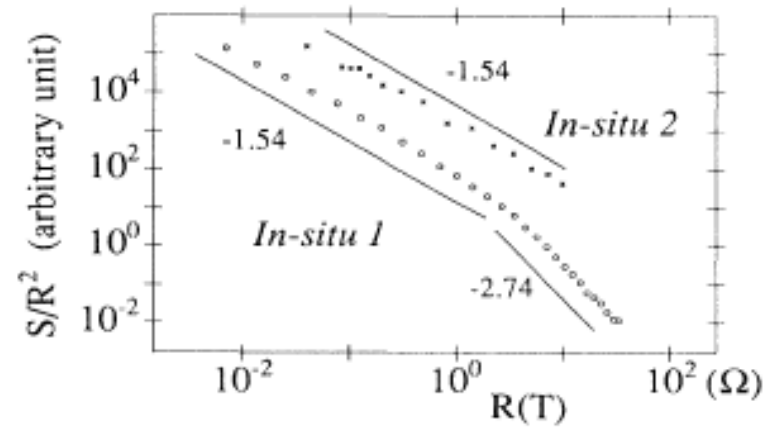


FIG. 1. Normalized noise versus resistance in high-quality high- T_c superconducting thin films in the percolation region. *In situ 1*: sample fabricated by coevaporation. *In situ 2*: sample fabricated by dc magnetron sputtering. The solid lines correspond to slopes predicted by scaling theory.

Ultra-low noise temperature control (originally developed by Per Nordblad, which we modified for the new needs)



- **P. Nordblad**, “Magnetic Anisotropy and Magnetic Phase Transitions of Iron and Manganese Compounds”, in Abstracts of Uppsala Dissertations from the Faculty of Science 556, Acta Universitatis Upsaliensis (Uppsala, Sweden, 1980).

- P. Svedlindh, K. Gunnarsson, P. Nordblad, L. Lundgren, H. Aruga, and A. Ito, Phys. Rev. B 40, 7162 (1989).

- ***Copper thermometer, (DC-) heater and the film sample on the same copper block;***

- ***in vacuum, thus passive thermal relaxation time > 1000 seconds;***

- ***thermometer with 4-point driving/probing arrangement with a differential transformer;***

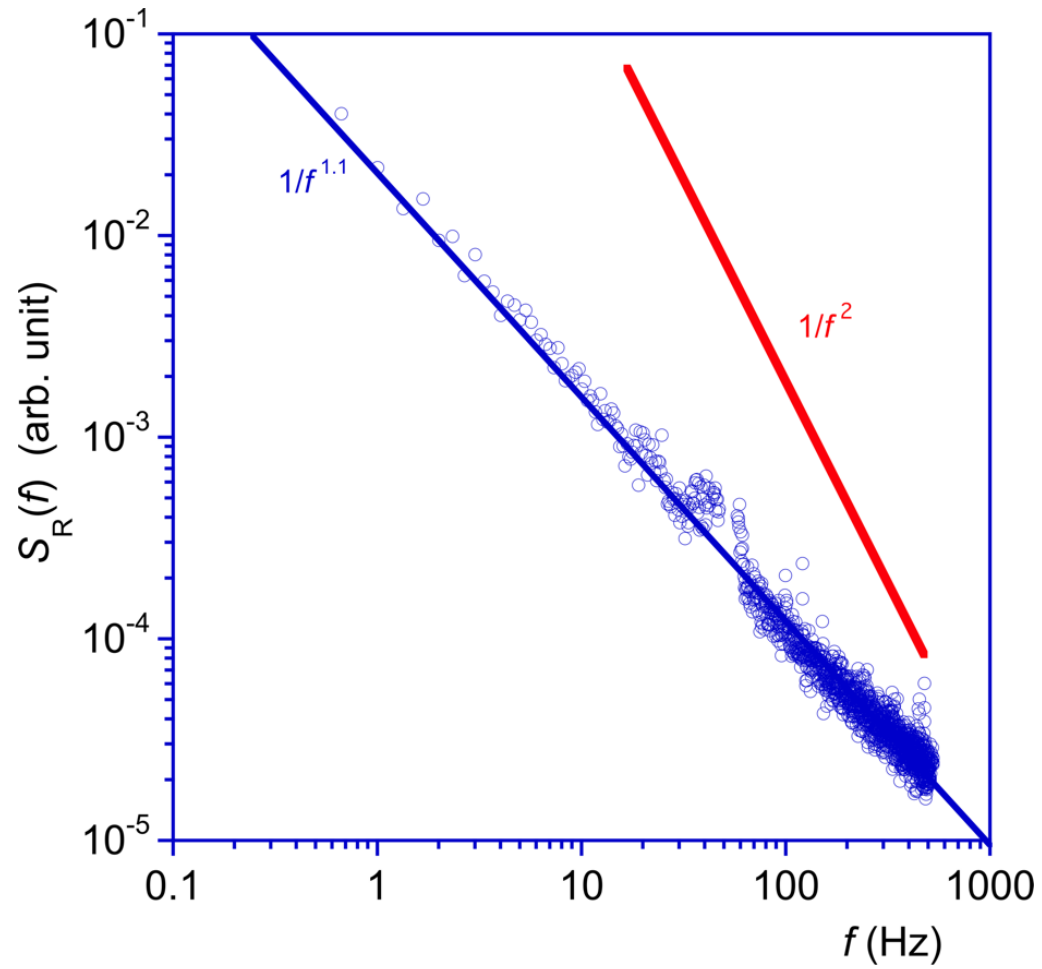
- ***temperature measurement bridge is of high-stability resistors in oil bath***

- ***which is driven by AC, 473Hz to reduce 50 Hz harmonics, generated, filtered/measured by a lockin amplifier***

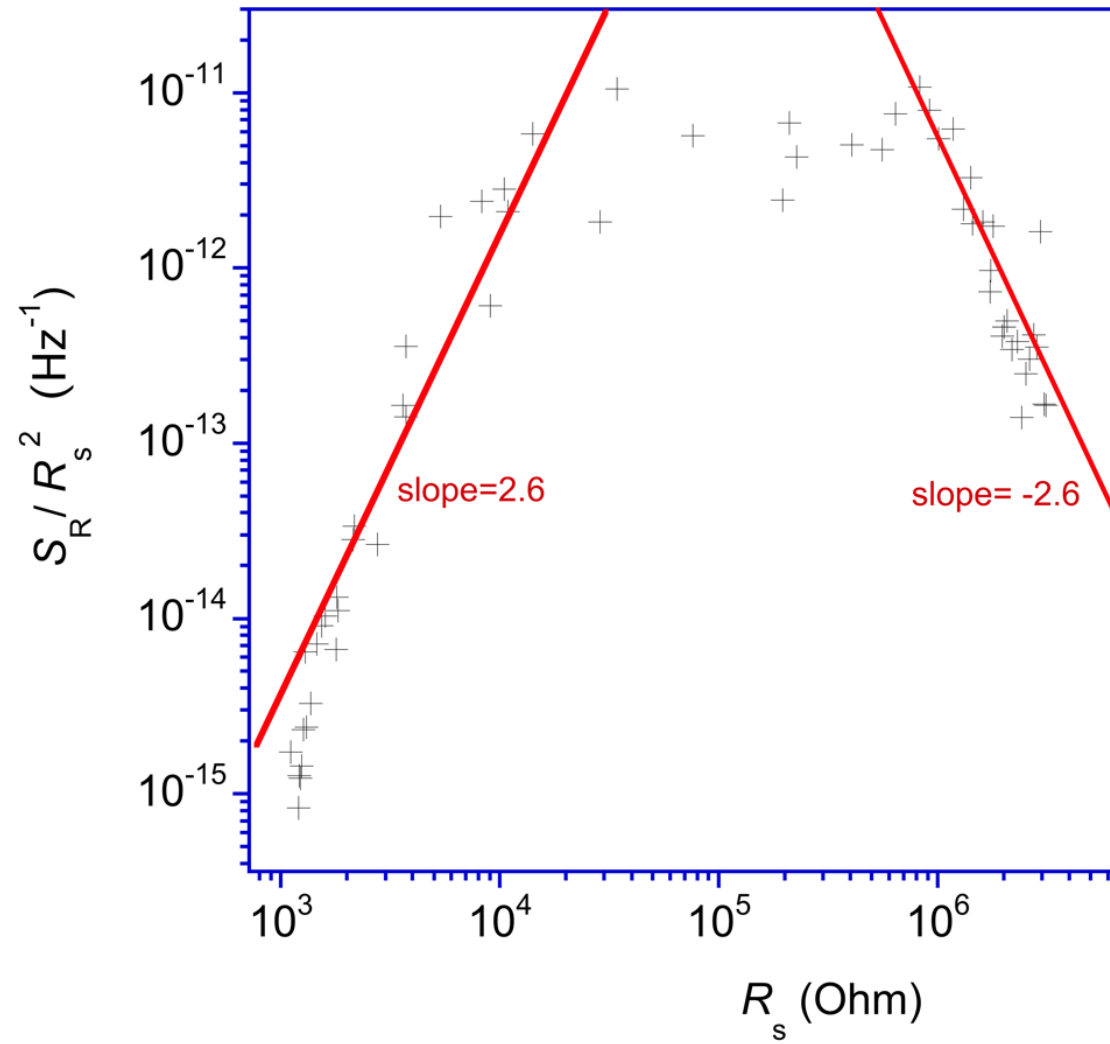
- ***lockin DC output is driving a PID controller, which drives the analog DC heater amplifier.***

Temperature noise less than 10^{-9} K/Hz^{0.5} can be achieved.

Conductance noise spectrum and checking for temperature fluctuations ($1/f^2$)



Scaling plot of the normalized noise versus the resistance



Our measured/fitted noise exponents in various high- T_c superconductor films (1989-1994)

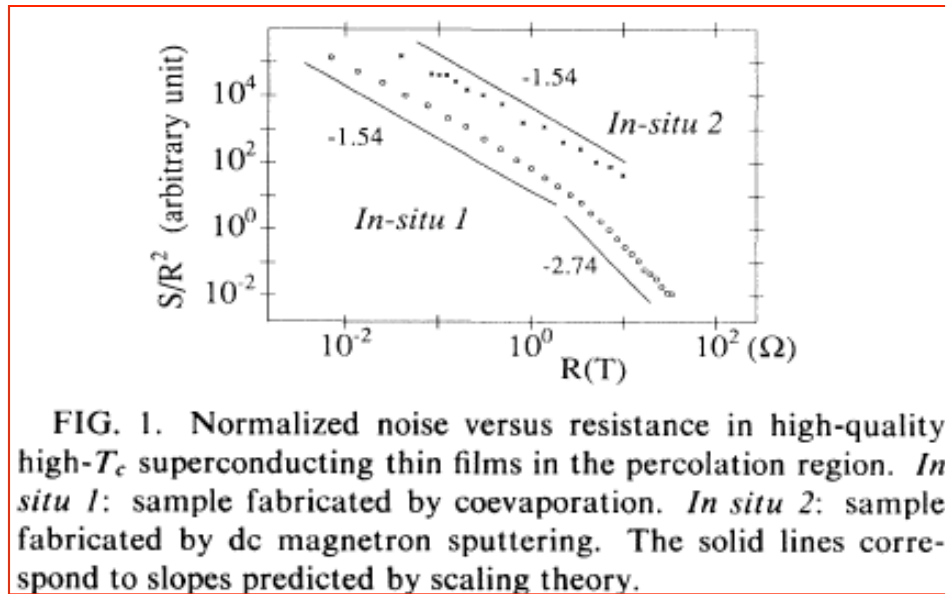
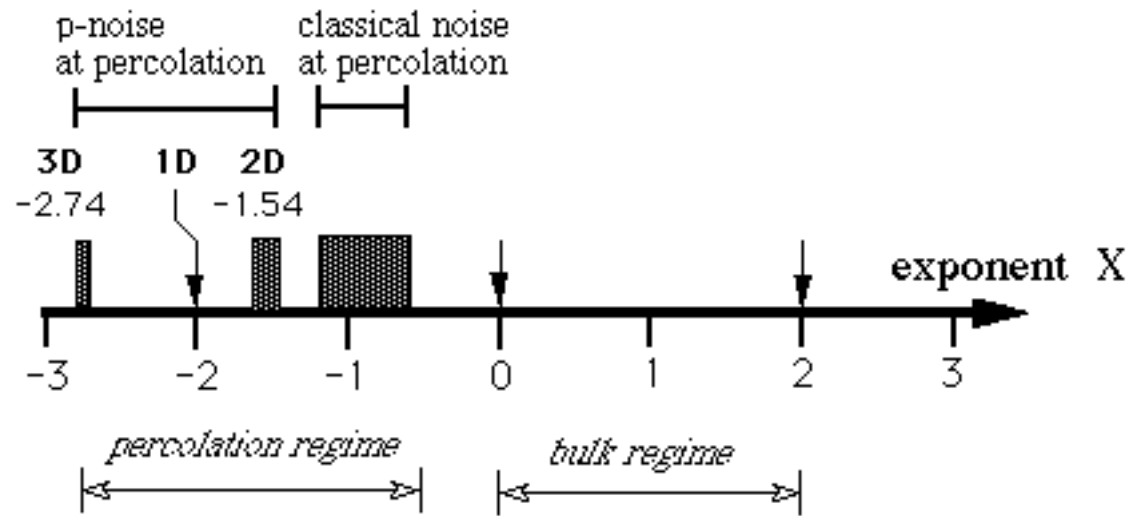
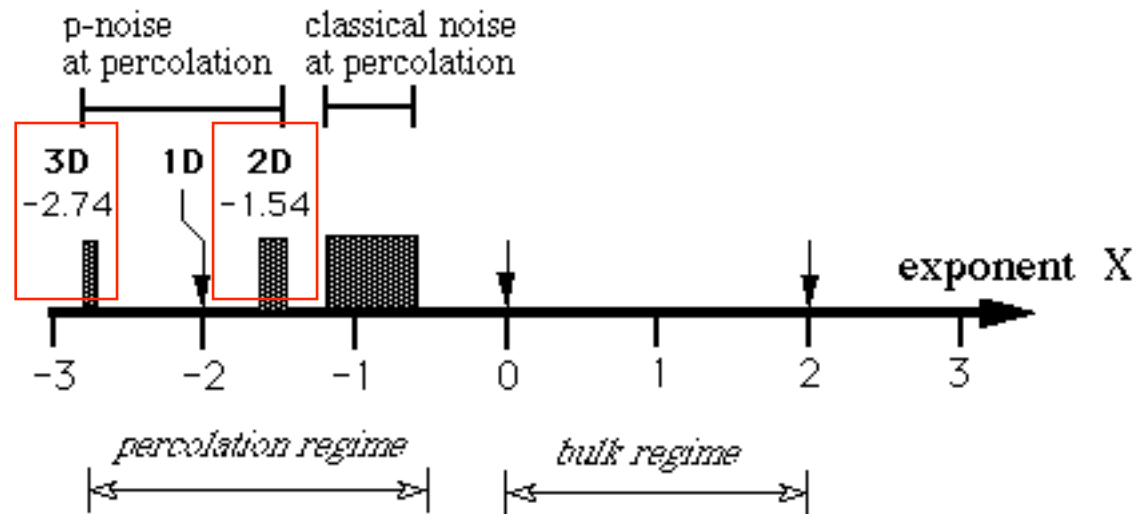


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Our measured/fitted noise exponents in various high- T_c superconductor films (1989-1994)



The -2.7 exponent would be fine at the high-temperature end however that is the 3D case where duality [P.M. Hui and D. Stroud, Phys. Rev. B 34, 8101 (1986)] does not force the same absolute value in at the low-temperature end as at the high one.

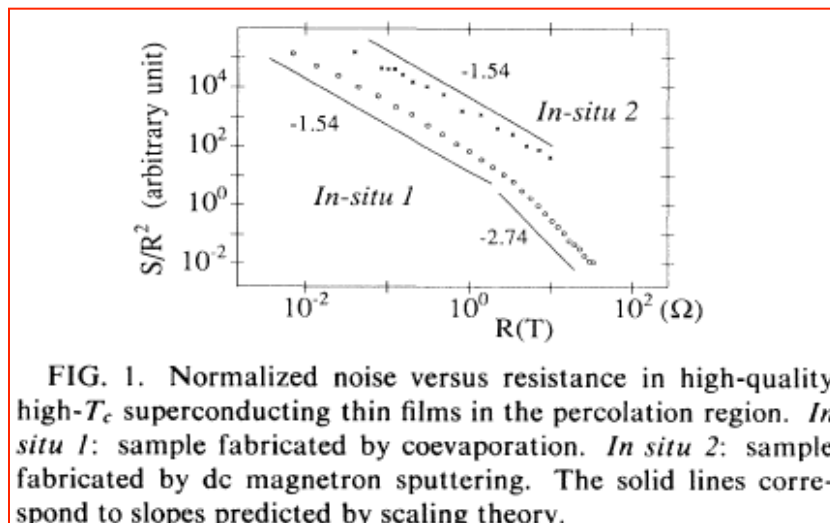


FIG. 1. Normalized noise versus resistance in high-quality high- T_c superconducting thin films in the percolation region. *In situ 1*: sample fabricated by coevaporation. *In situ 2*: sample fabricated by dc magnetron sputtering. The solid lines correspond to slopes predicted by scaling theory.

Kiss-Svedlindh, PRL 1993.

TABLE I. Scaling exponents of the resistance and the normalized noise in random resistor networks. The resistance of the normal-conducting elements fluctuates independently; i.e., the number of noise sources is given by the number of resistors.

	RS composite ($p_r > p_c$)		RI composite ($p_r > p_c$)	
	$R_{rs} \propto (p_r - p_c)^s;$	$\frac{S_{rs}(f)}{R_{rs}^2} \propto R_{rs}^{-l_{rs}}$	$R_{ri} \propto (p_r - p_c)^{-t};$	$\frac{S_{ri}(f)}{R_{ri}^2} \propto R_{ri}^{-l_{ri}}$
	s	l_{rs}	t	l_{ri}
1D	1	1
2D	1.297 ± 0.07	0.86 ± 0.02	1.297 ± 0.07	0.86 ± 0.02
3D	0.73 ± 0.011	0.9 ± 0.32	1.96 ± 0.1	0.80 ± 0.1

[Kiss-Svedlindh, p-fluctuations, percolation noise model, 1993](#)

TABLE II. Comparison of classical (l) and new (λ) noise exponents.

	RS composite ($p_r > p_c$)		RI composite ($p_r > p_c$)	
	$\frac{S_{rs}(f)}{R_{rs}^2} \propto R^{-l_{rs}};$	$\frac{S_{rs}(f)}{R_{rs}^2} \propto R_{rs}^{-\lambda_{rs}}$	$\frac{S_{ri}(f)}{R_{ri}^2} \propto R_{ri}^{-l_{ri}};$	$\frac{S_{ri}(f)}{R_{ri}^2} \propto R_{ri}^{-\lambda_{ri}}$
	l_{rs}	$\lambda_{rs} = 2/s$	l_{ri}	$\lambda_{ri} = 2/t$
1D	1	2
2D	0.86 ± 0.02	1.54 ± 0.09	0.86 ± 0.02	1.54 ± 0.09
3D	0.9 ± 0.32	2.74 ± 0.04	0.80 ± 0.1	1.02 ± 0.05

Scaling Law of Resistance Fluctuations in Stationary Random Resistor Networks

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(Received 24 July 2000)

In a random resistor network we consider the simultaneous evolution of two competing random processes consisting in breaking and recovering the elementary resistors with probabilities W_D and W_R . The condition $W_R > W_D/(1 + W_D)$ leads to a stationary state, while in the opposite case, the broken resistor fraction reaches the percolation threshold p_c . We study the resistance noise of this system under stationary conditions by Monte Carlo simulations. The variance of resistance fluctuations $\langle \delta R^2 \rangle$ is found to follow a scaling law $|p - p_c|^{-\kappa_0}$ with $\kappa_0 = 5.5$. The proposed model relates quantitatively the defectiveness of a disordered media with its electrical and excess-noise characteristics.



The **Pennetta-Trefan-Reggiani** model of "dynamical percolation" with microscopic damage and healing processes with separate rates in 2D produces 2.6 exponents in the steady-state at low-temperature.

In 2D, *due to duality* [P. M. Hui and D. Stroud, Phys. Rev. B 34, 8101 (1986)], the same absolute exponent value holds in the high-temperature limit with negative sign.

tain, in the disordered network regime,

$$\frac{\langle \delta R^2 \rangle}{\langle R \rangle^2} \sim \langle R \rangle^{-s} \quad (8)$$

with $s = 2.6$ as reported in Fig. 6.

in the percolation (scaling) region:

$$\frac{S_R(f)}{R^2} \propto R^x \quad \text{with} \quad x \approx \pm 2.6$$

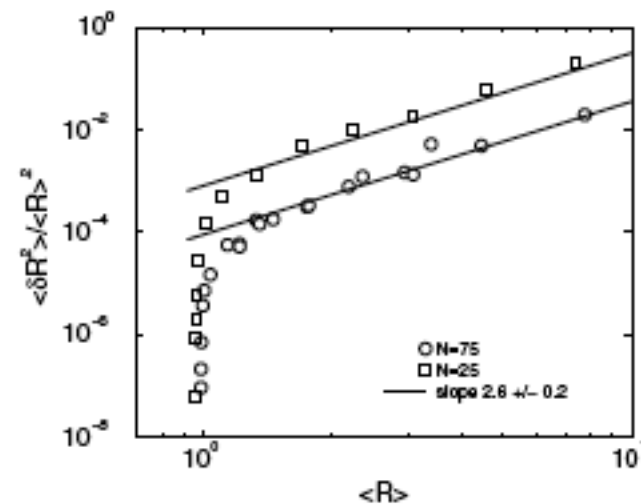
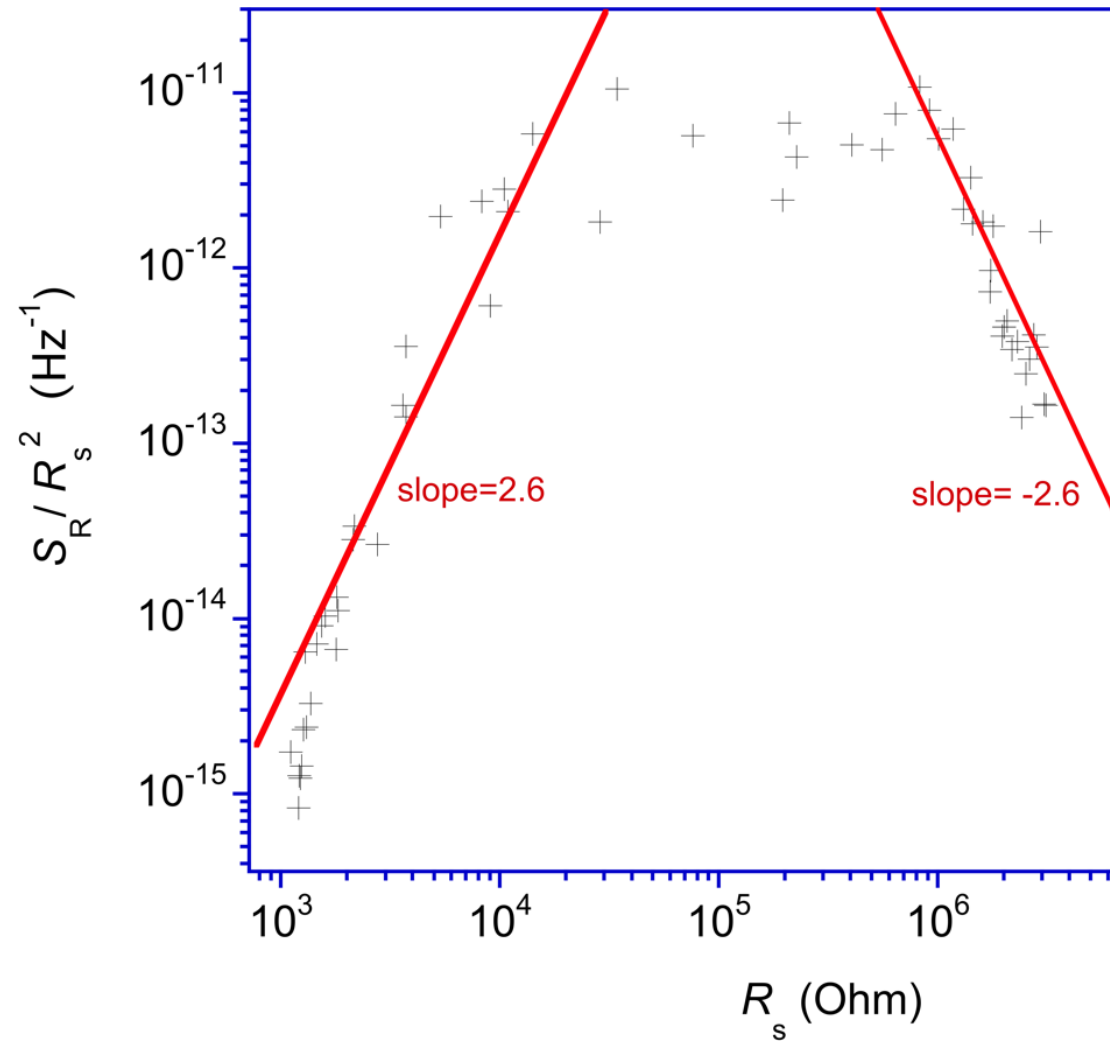


FIG. 6. Resistance noise normalized to the square of the average RRN resistance $\langle \delta R^2 \rangle / \langle R \rangle^2$ as a function of the average resistance $\langle R \rangle$.

Scaling plot of the normalized noise versus the resistance



UPoN!!!

- In the PTR model, the noise is not inherent in the resistance but comes from the switching
- PTR see Lorentzian spectra. How do we get $1/f$ noise?
- Hierarchy of switching time constants?
- Why is the spectrum $1/f$ in the whole temperature range?
- Perhaps another model is relevant?