



On the WEAK MEASUREMENT
of the ELECTRICAL THz CURRENT:
a NEW SOURCE of NOISE?

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① Open problem

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- ② Novel approach to model the measurement of THz current

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How do we model the measurement of the high frequency current?

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Year	2015	2020	2025
Cutoff Frequency (GHz)	620	1137	2062

International Technology Roadmap for Semiconductors (2011)

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Continuous (or very High Frequency) Measurement



Inclusion of the Back-Action (Disturbance of quantum system)

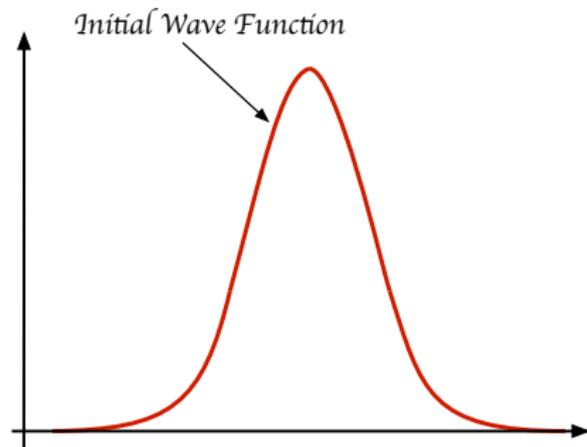
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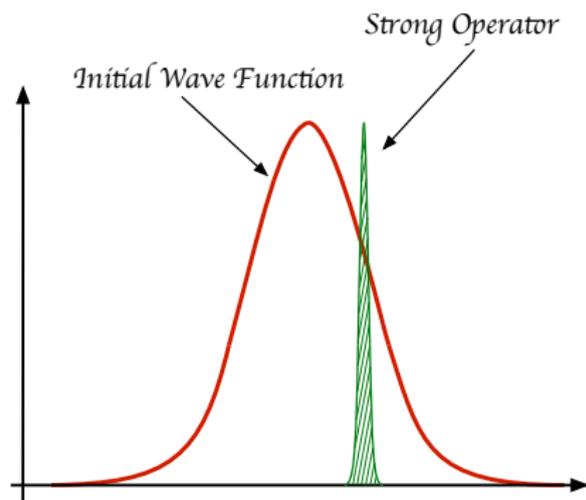


J. Von Neumann: PUP (1955)

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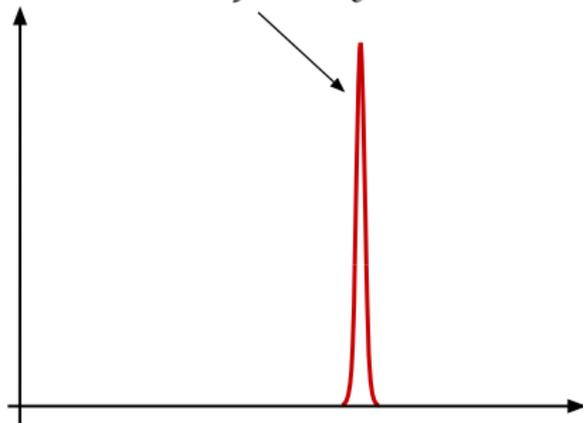
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Wave Function after Strong Measurement

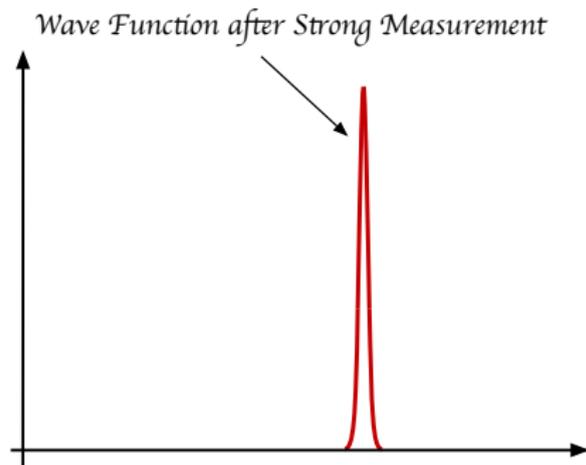


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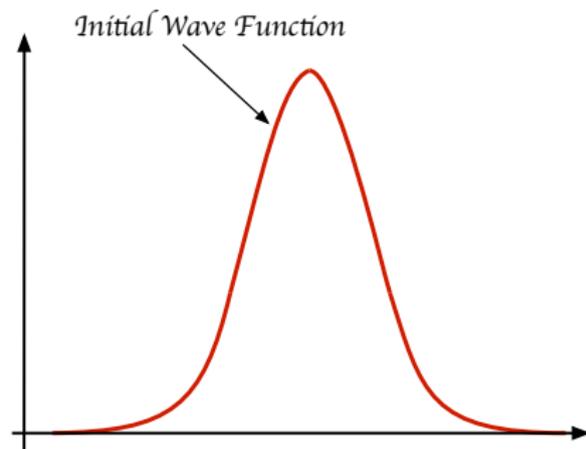
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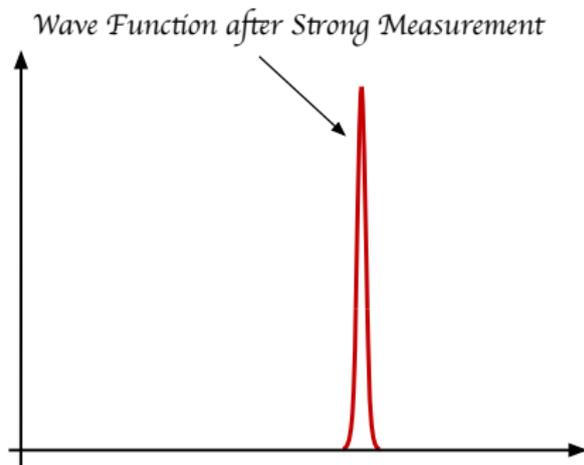


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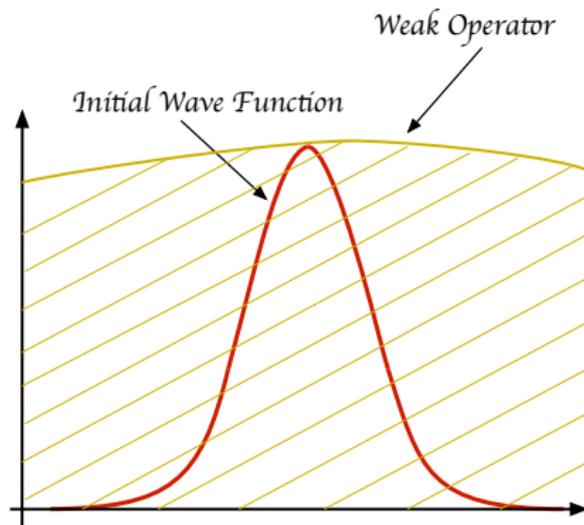
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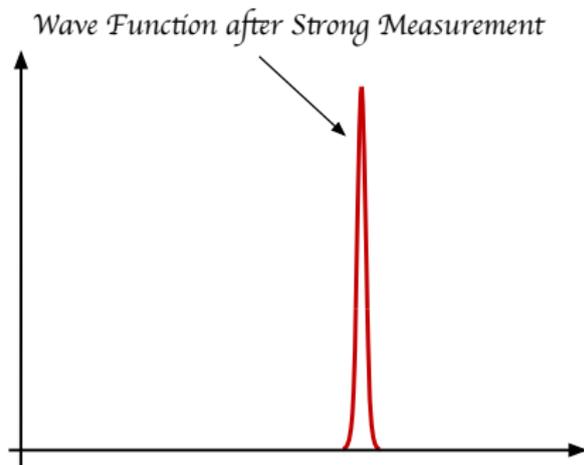


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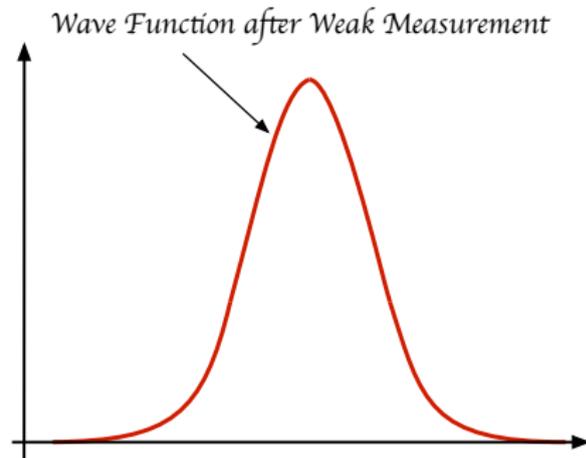
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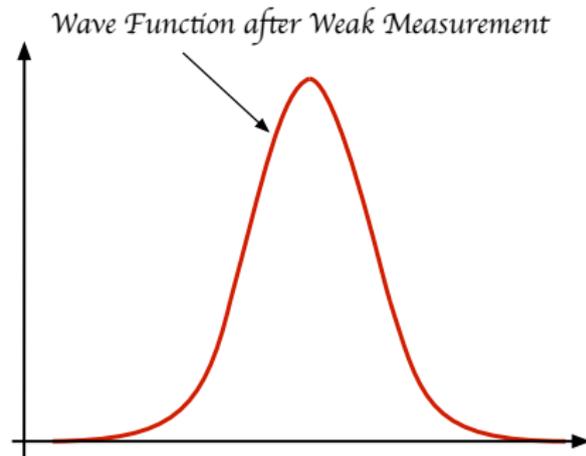
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How is Measurement modelled in Quantum Mechanics?

Weak

Main features of the *Weak* Measurement

- $\langle I \rangle_{strong} = \langle I \rangle_{weak}$
- Wave Function of the system is slightly perturbed after the interaction



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Measuring the current at high frequency

Three ways:

- | Taking information from the system without worrying about the apparatus → **Be careful!**

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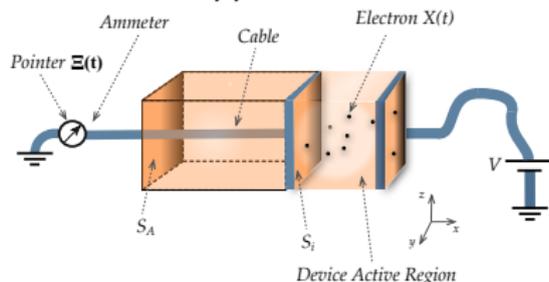
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 - *Include the apparatus and see what happens!*

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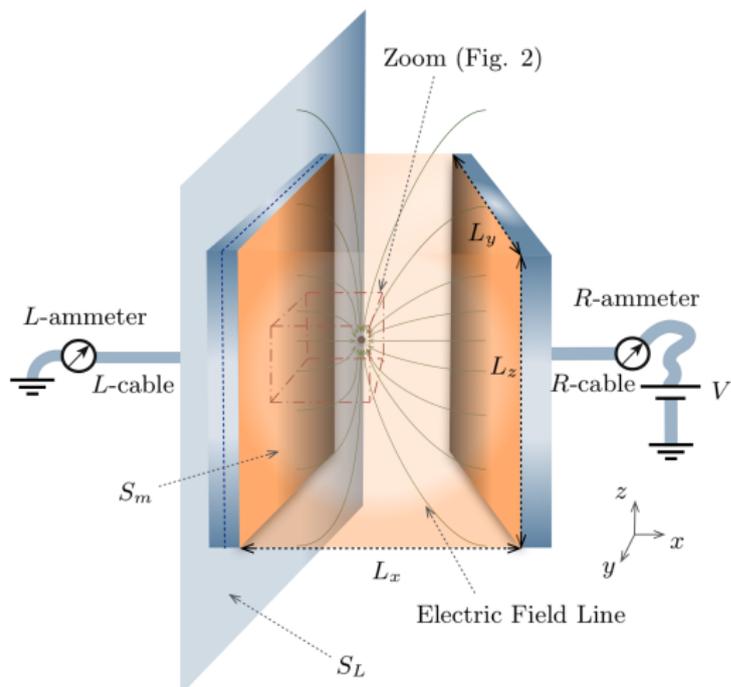


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2 - Approach to model the measurement of THz current

Many-Body Problem

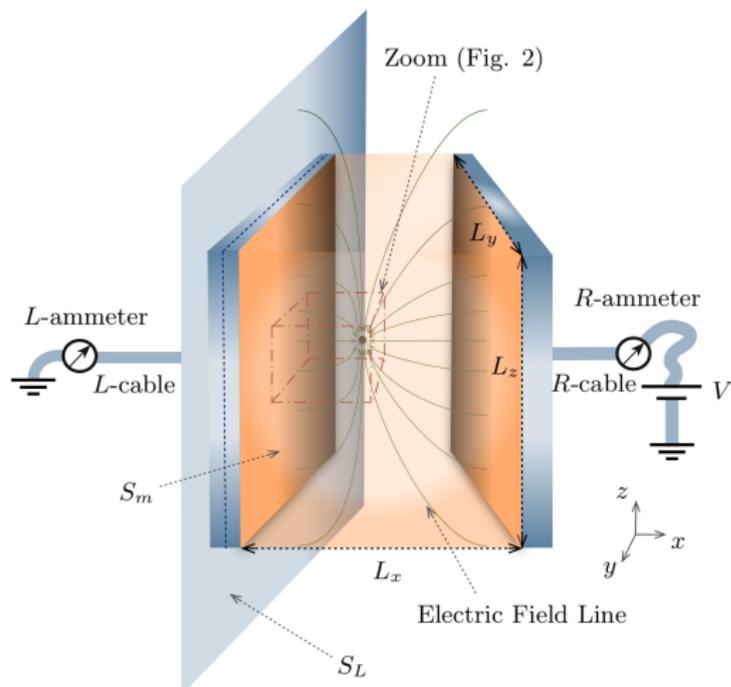
$$i\hbar \frac{\partial \Psi(x_1, x_2, \dots, x_N, t)}{\partial t} = [H_0 + V_{Coul}] \Psi(x_1, x_2, \dots, x_N, t)$$



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• Coulomb Interaction

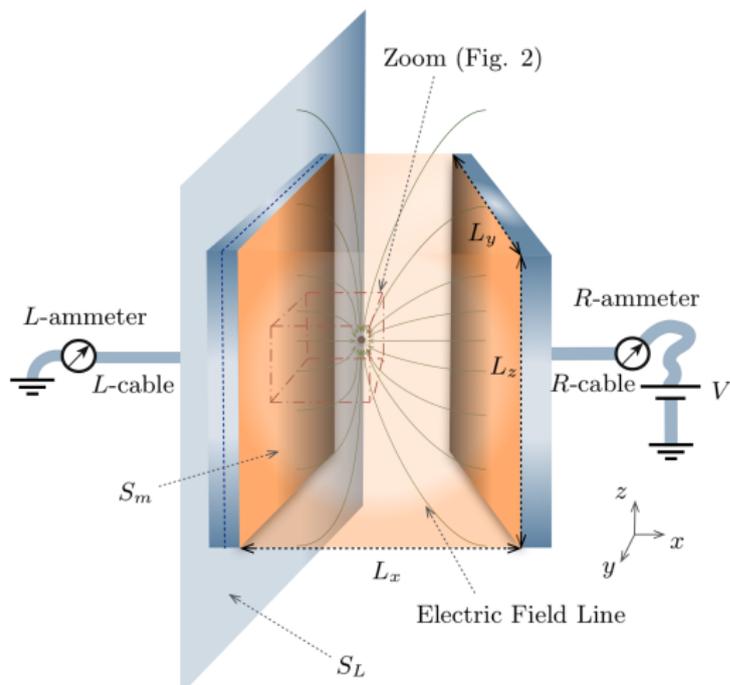
↓

$$V_{Coul} = \frac{1}{4\pi\epsilon(\mathbf{r})} \frac{1}{2} \sum_{i=1, j \neq i}^N \frac{q_i q_j}{\sqrt{(x_i - x_j)^2}}$$

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- Many Particle

↓

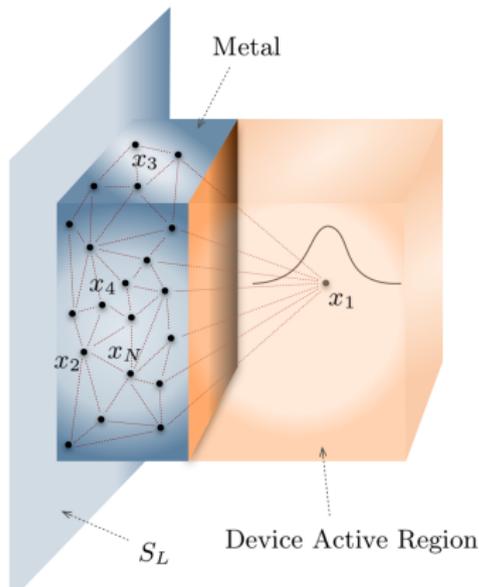
Numerically inaccessible

2 - Approach to model the measurement of THz current

Conditional Wave Function

$$i\hbar \frac{\partial \psi(x_1, t)}{\partial t} = [H_0 + V_{Cond}] \psi(x_1, t) \quad \frac{dX_1(t)}{dt} = \frac{\hbar}{m} \text{Im} \left(\frac{\nabla \psi}{\psi} \right)$$

D. Dürr, S. Goldstein, and N. Zanghi: JSP (1992;2005)

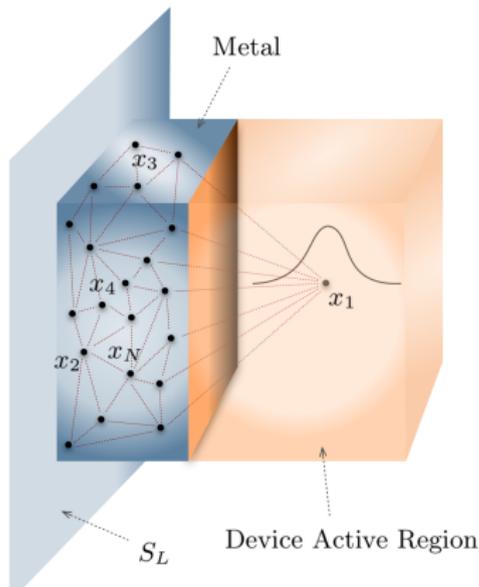


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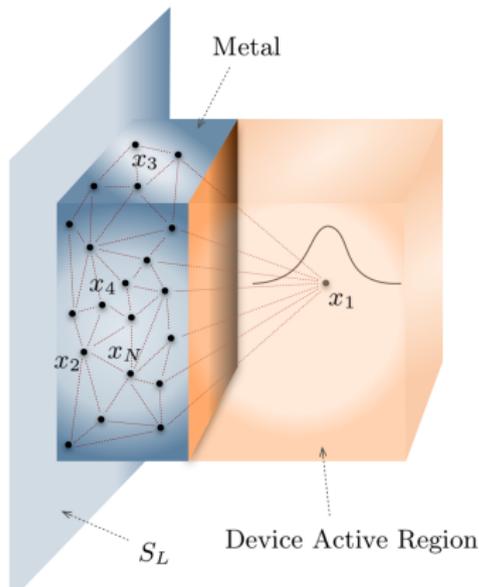
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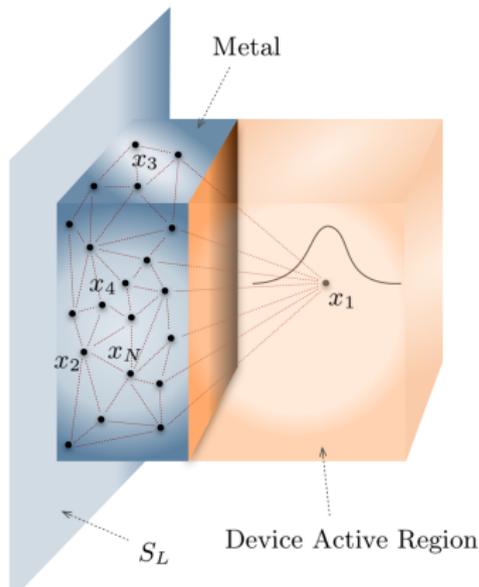
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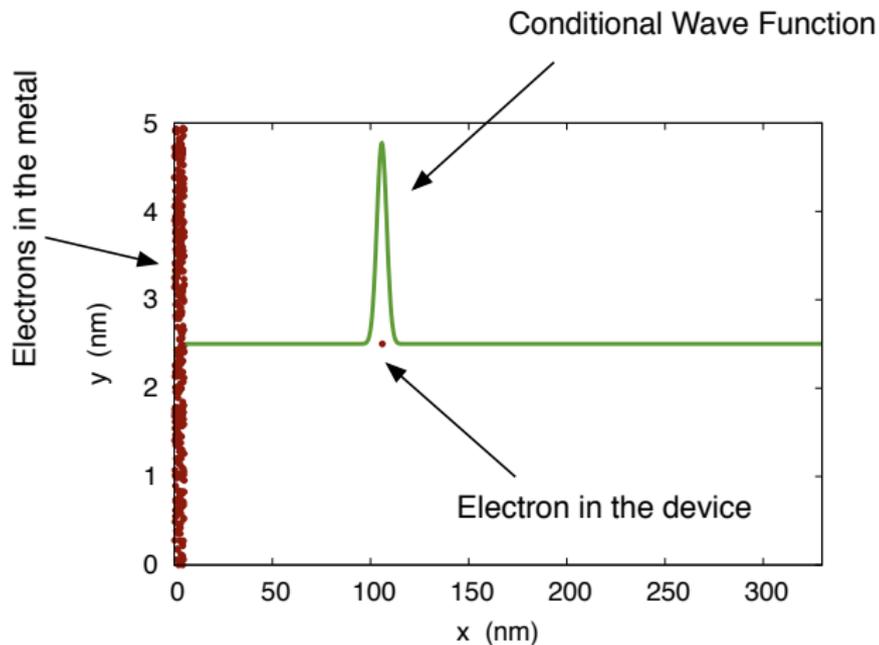
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- $I_d(t) = \int_{S_L} \epsilon(r) \frac{dE(r; X_1(t), \dots, X_N(t), t)}{dt} \cdot ds$

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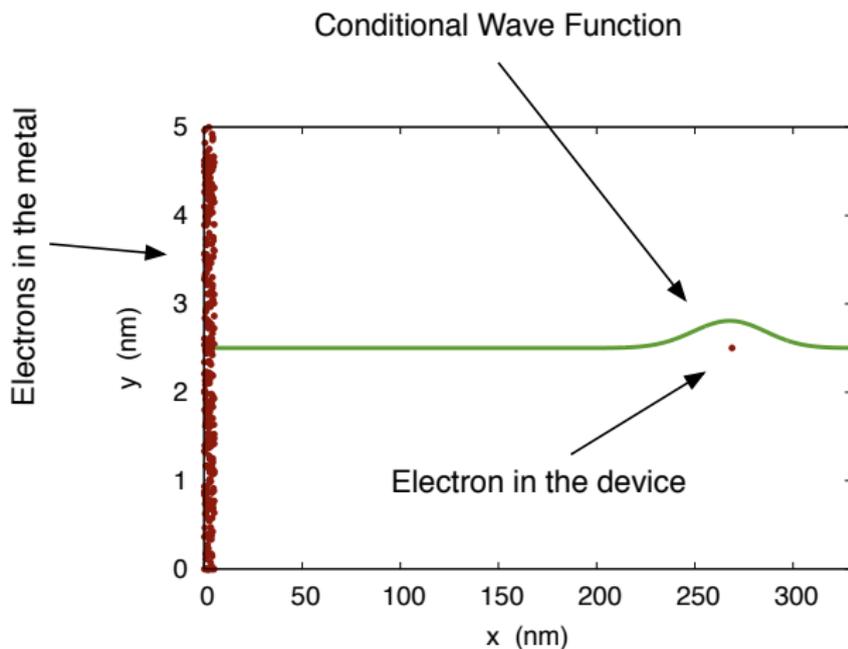


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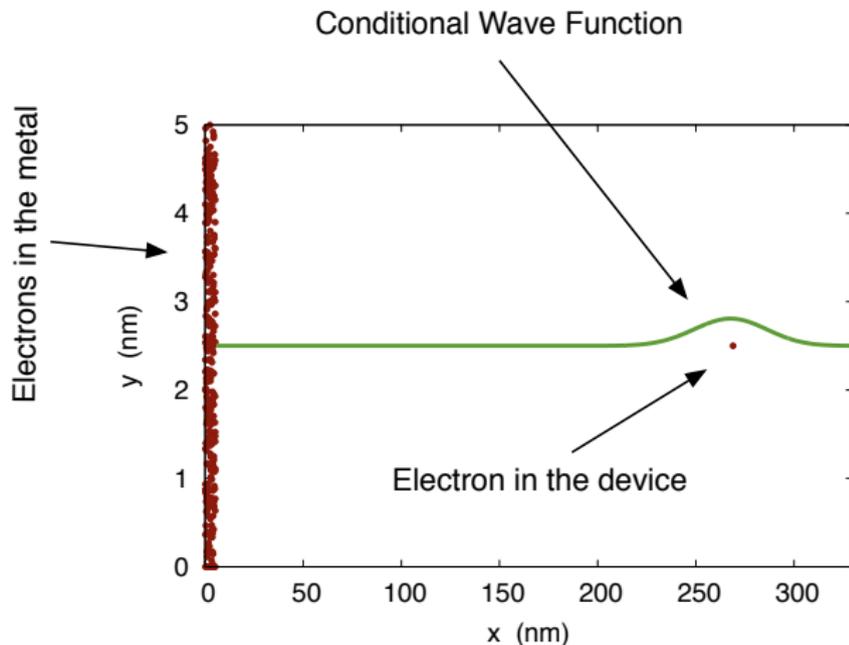
Evolution of the system under the interaction with the apparatus

“

2 - Approach to model the measurement of THz current



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We have a numerical method to tackle the many-body problem!

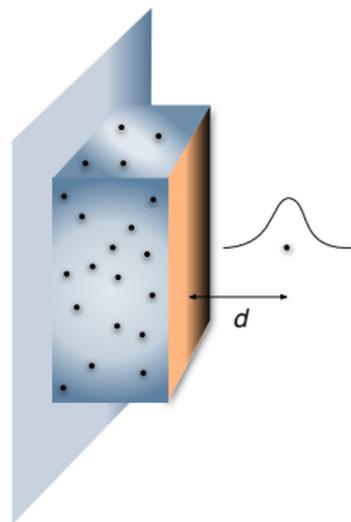


We include the **Back-Action** of the measuring apparatus!!!

- ① Open problem
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- ③ **New source of noise**
- ④ Measurement of the local (Bohmian) velocities
- ⑤ Concluding Remarks

3 - New source of noise

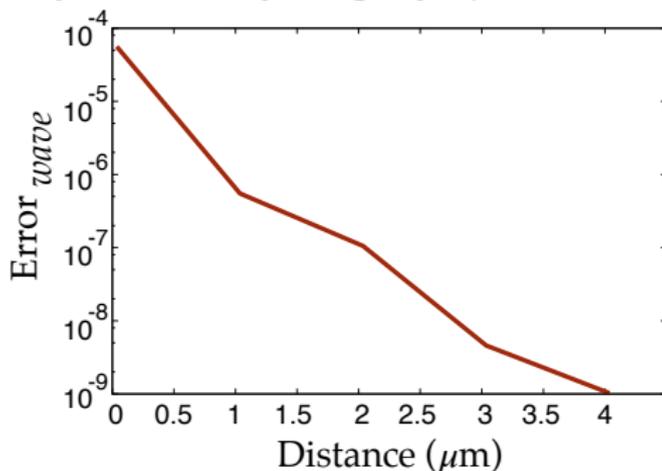
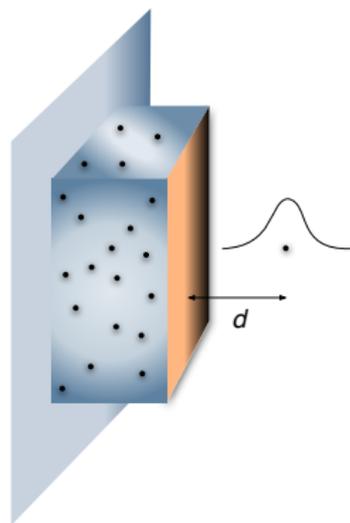
The wave function of the system is only “slightly” perturbed



Different distances from the metal
surface \rightarrow Parameter d

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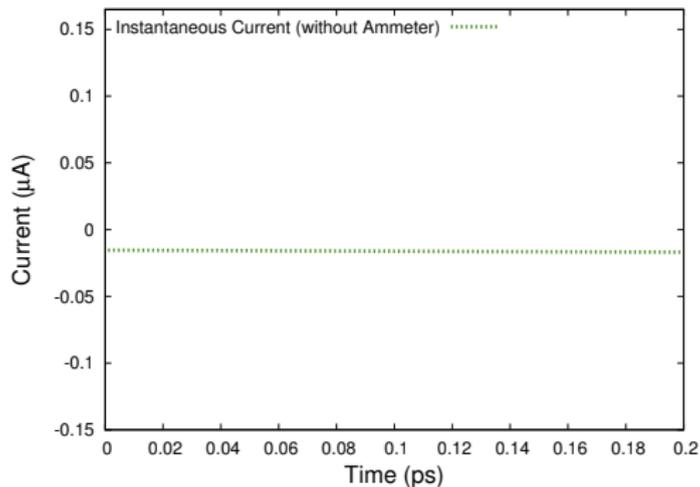
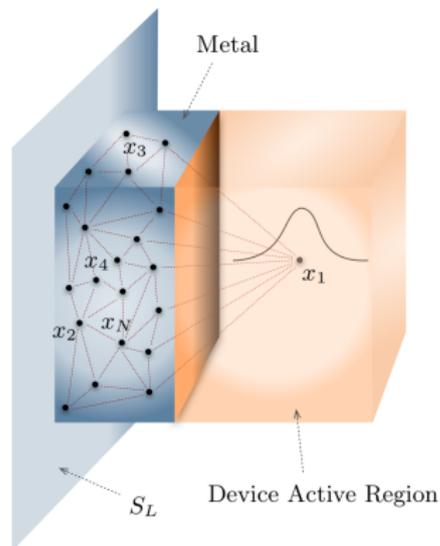
Different distances from the metal surface → Parameter d

Error Wave Function
$$\text{Error}_{\text{Wave}} = \int |\psi_{\text{int}} - \psi_{\text{free}}|^2 dx$$

The error in the Wave Function decreases with the distance!

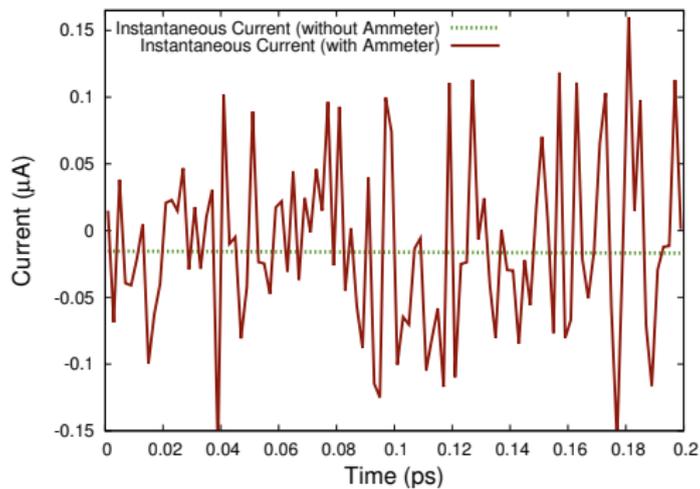
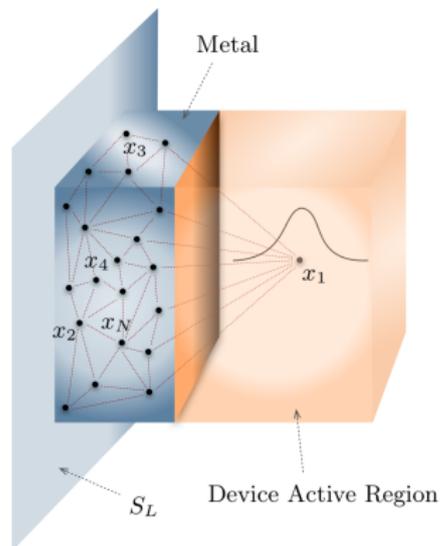
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Total Current Measured in the Surface S_L



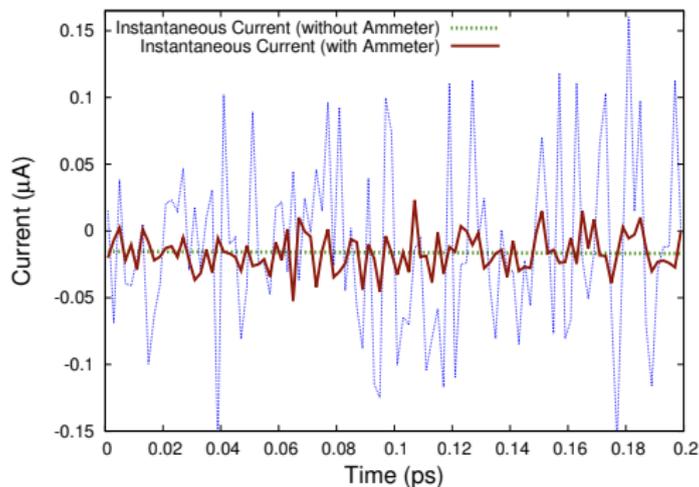
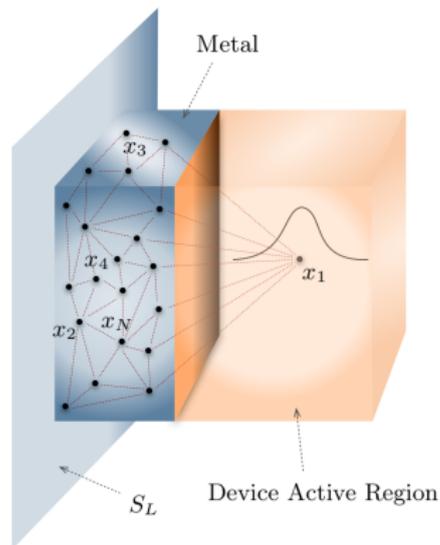
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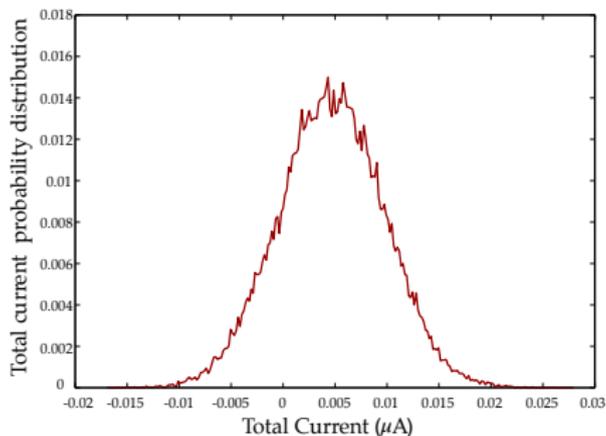


Repeating many times the experiment

we reach the mean value!!!

3 - New source of noise

Probability distribution of the total current in a large metallic surface



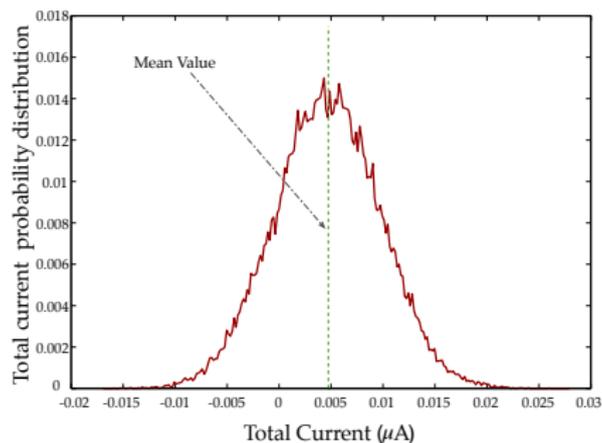
Output current very noisy!



Additional source of noise!!!

3 - New source of noise

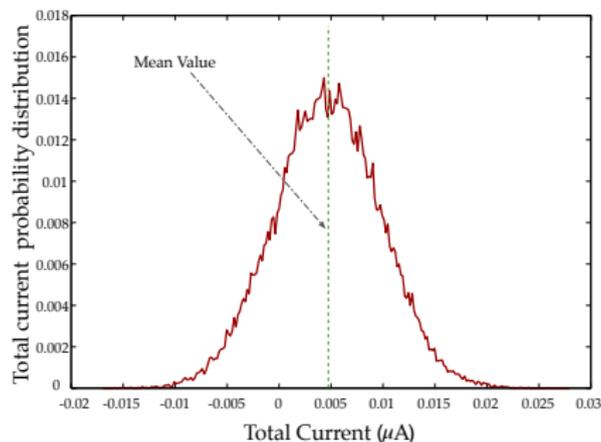
Probability distribution of the total current in a large metallic surface



$$\langle I \rangle_{strong} = \langle I \rangle_{weak}$$

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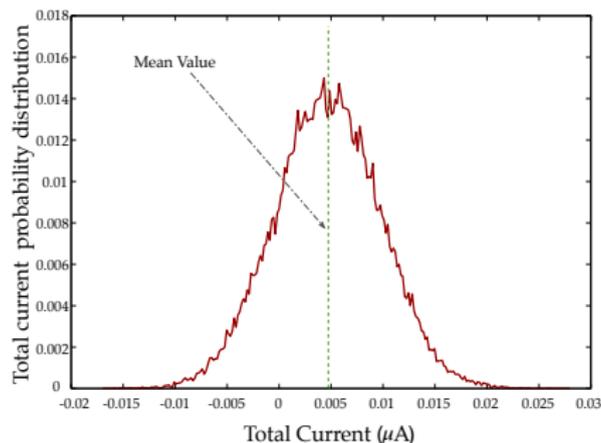
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⇒ The error in the Wave Function decreases with the distance!

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Weak Measurement!!!

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4 - Measurement of the local (Bohmian) velocities

Total Current (conduction *plus* displacement components)

$$\langle I(t) \rangle = \int_{S_i} \langle \mathbf{J}_c(\mathbf{r}, t) \rangle \cdot d\mathbf{s} + \int_{S_i} \epsilon \frac{d\langle \mathbf{E}(\mathbf{r}, t) \rangle}{dt} \cdot d\mathbf{s}$$

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Or from the Ramo-Schockley-Pellegrini theorem

$$\langle I(t) \rangle = -\int_{\Omega} \mathbf{F}(\mathbf{r}) \cdot \langle \mathbf{J}_c(\mathbf{r}, t) \rangle \cdot d\mathbf{v} + \int_S \epsilon \cdot \mathbf{F}(\mathbf{r}) \cdot \frac{d\langle V(\mathbf{r}, t) \rangle}{dt} \cdot d\mathbf{s}$$

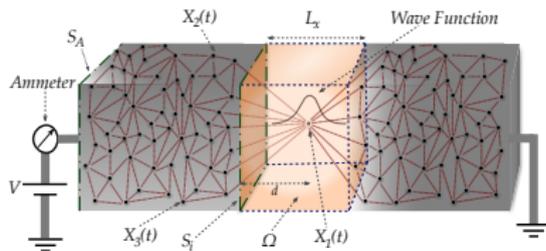
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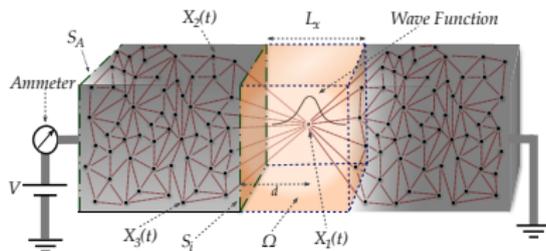
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In the case of two metallic surfaces S_i at distance $S_i \gg L_x^2$:

$$\langle I(t) \rangle = \frac{q}{mL_x} \langle p(t) \rangle_{\Omega}$$

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Total Current in a large surface



Weak Measurement of the Momentum!!!

4 - Measurement of the local (Bohmian) velocities

Weak measurement: a new way for measuring incompatible observables

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Measurement of position and momentum (for an ensemble of experiments)

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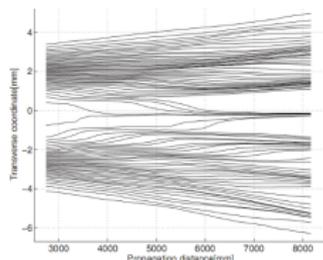
Weak measurement: a new way for measuring incompatible observables



Measurement of position and momentum (for an ensemble of experiments)



Experimental measurement of the local velocity



From: S. Kocis, *et al.*, *Science*, **332** (2011).

Experimental Bohmian trajectories:
Photons in a double slit set-up

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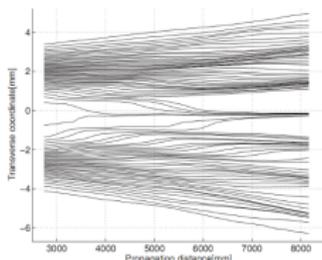
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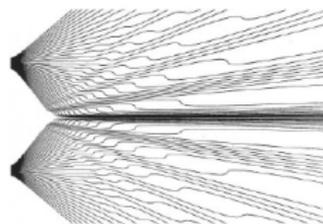


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Theoretical Bohmian trajectories

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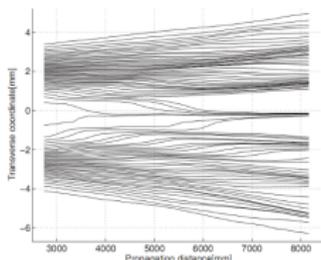
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Measurement of position and momentum (for an ensemble of experiments)

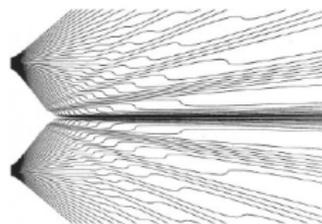


Experimental measurement of the local velocity



From: S. Kocis, *et al.*, *Science*, **332** (2011).

Experimental Bohmian trajectories:
Photons in a double slit set-up



From: C. Philippidis, *et al.*, *Il nuovo cimento B* (1979)
Theoretical Bohmian trajectories



Is it possible to envisage an analogous experiment for electrons?

4 - Measurement of the local (Bohmian) velocities

Experimental proposal for measuring the local (Bohmian) velocity

$$v(x) = \frac{1}{m} \operatorname{Re} \frac{\langle x | \hat{p} | \psi \rangle}{\langle x | \psi \rangle}$$

$$E[p_w | x_s] = \frac{\int dp_w p_w \mathcal{P}(p_w \cap x_s)}{\mathcal{P}(x_s)} = v(x_s, \tau)$$

H M Wiseman 2007 New J. Phys. 9 165

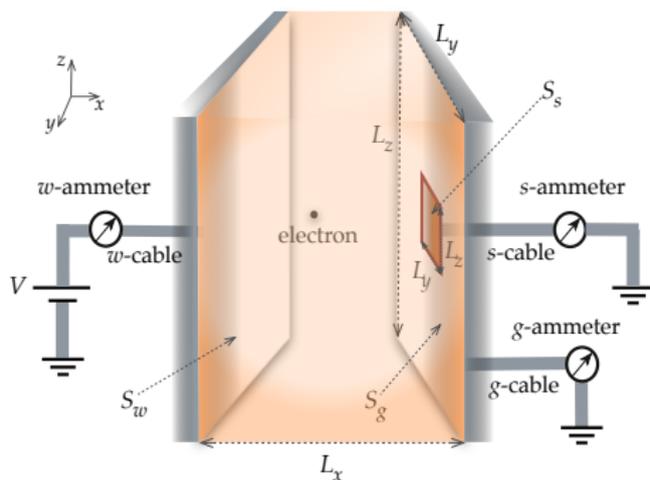
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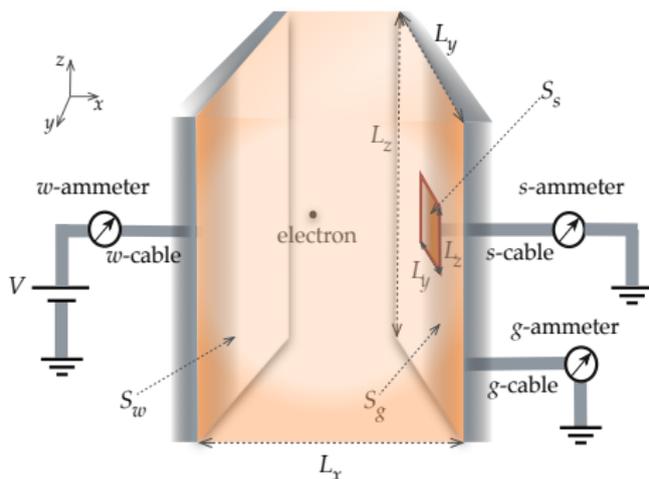
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- S_w where $L_y, L_z \gg L_x$
 $\rightarrow \langle I_w \rangle \propto \langle p_x \rangle$

WM of total current = WM of momentum

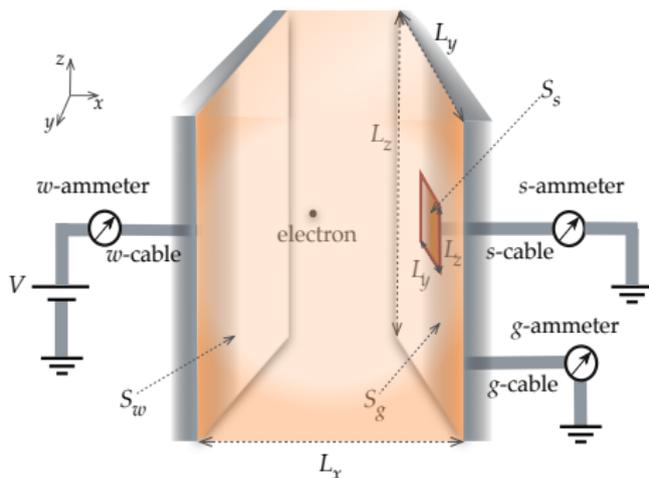
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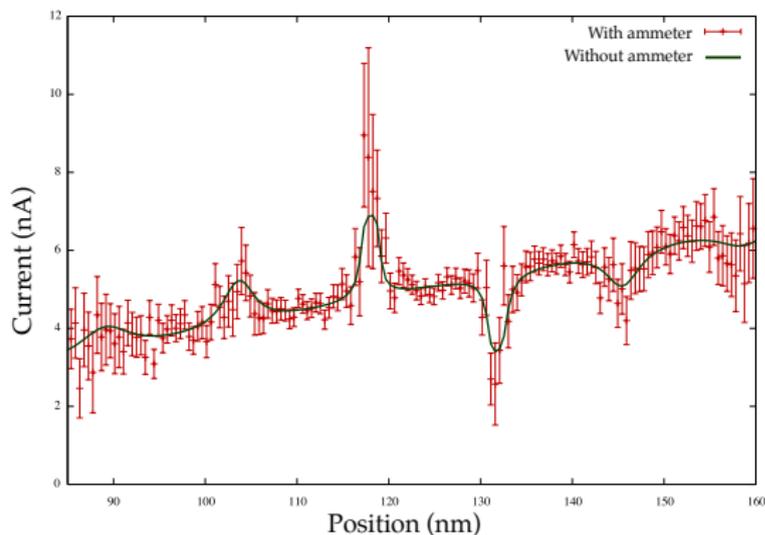
WM of total current = WM of momentum

- S_s where $L'_y, L'_z \ll L_x$
 $\rightarrow \langle I_s \rangle \propto |\langle r_s | \psi \rangle|^2$

Post-selection with position measurement

4 - Measurement of the local (Bohmian) velocities

Numerical Experiments



$$\begin{aligned} E[p_w | x_s] &= \frac{\int dp_w p_w \mathcal{P}(p_w \cap x_s)}{\mathcal{P}(x_s)} = \\ &= \frac{J(x_s, \tau)}{|\psi(x_s, \tau)|^2} \equiv v(x_s, \tau) \end{aligned}$$

Weak Measurement

$$\langle I_w \rangle \propto \langle p_x \rangle$$

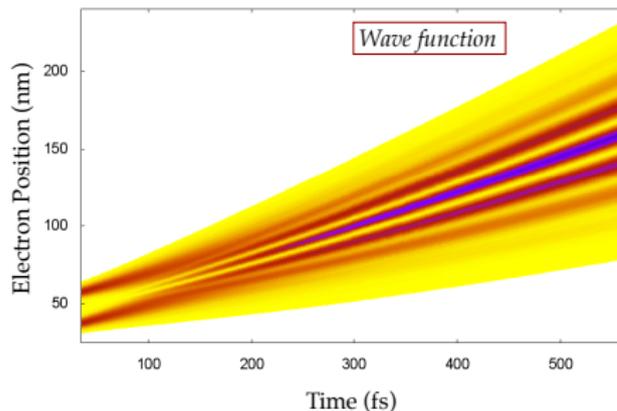
Strong Measurement

$$\langle I_s \rangle \propto |\langle r_s | \psi \rangle|^2$$

Reconstruction of the Bohmian velocity
from an ensemble of 55000
numerical experiments

4 - Measurement of the local (Bohmian) velocities

Numerical Experiments



Wave function

$$\begin{aligned} E[p_W | x_S] &= \frac{\int dp_W p_W \mathcal{P}(p_W \cap x_S)}{\mathcal{P}(x_S)} = \\ &= \frac{J(x_S, \tau)}{|\psi(x_S, \tau)|^2} \equiv v(x_S, \tau) \end{aligned}$$

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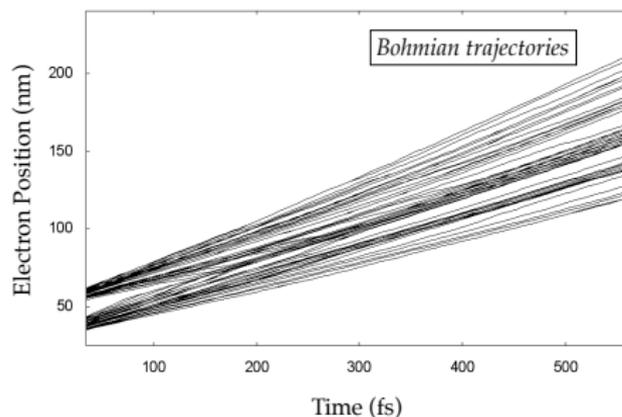
$$\langle I_W \rangle \propto \langle p_X \rangle$$

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Bohmian trajectories

$$E[p_W | X_S] = \frac{\int dp_W p_W \mathcal{P}(p_W \cap X_S)}{\mathcal{P}(X_S)} =$$
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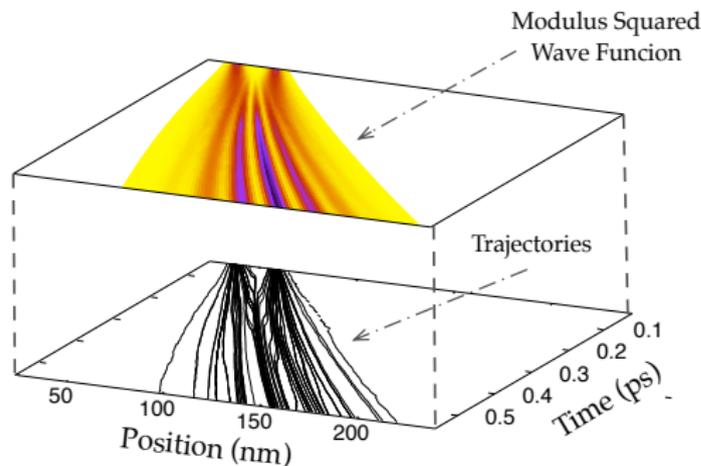
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Wave function *and* Bohmian trajectories

D. M., X. Oriols and N. Zanghì: *Weak Values from Displacement Currents in Multiterminal Electron Devices*, in preparation.
D. Dürr, S. Goldstein, and N. Zanghì, *On the Weak Measurement of Velocity in Bohmian Mechanics*, *Journal of Statistical Physics* 134, 1023 (2009).

- ① Open problem
- ② Novel approach to model the measurement of THz current
- ③ New source of noise
- ④ Measurement of the local (Bohmian) velocities
- ⑤ **Concluding Remarks**

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- Reconstruction of the Bohmian trajectories in a solid state device

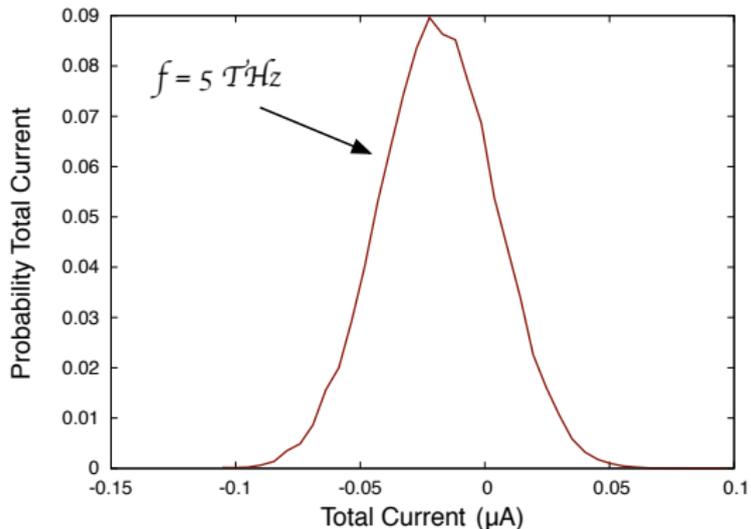


On the WEAK MEASUREMENT
of the ELECTRICAL THz CURRENT:
a NEW SOURCE of NOISE?

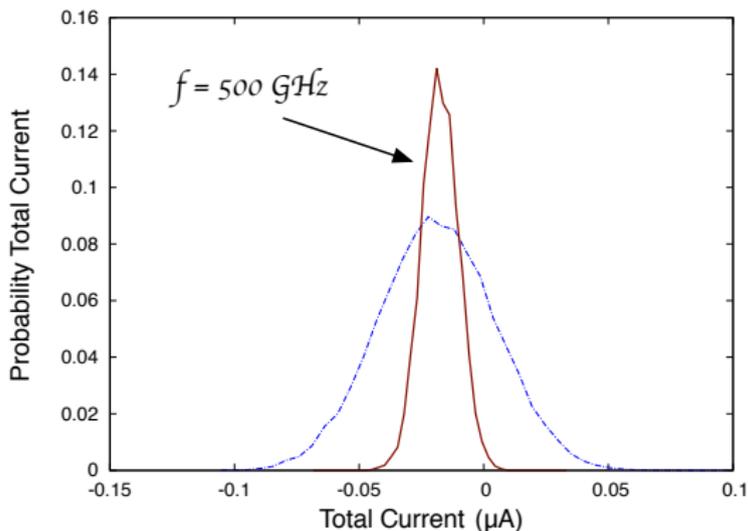
Damiano Marian, Nino Zanghì, Xavier Oriols

THANKS FOR THE ATTENTION!!!

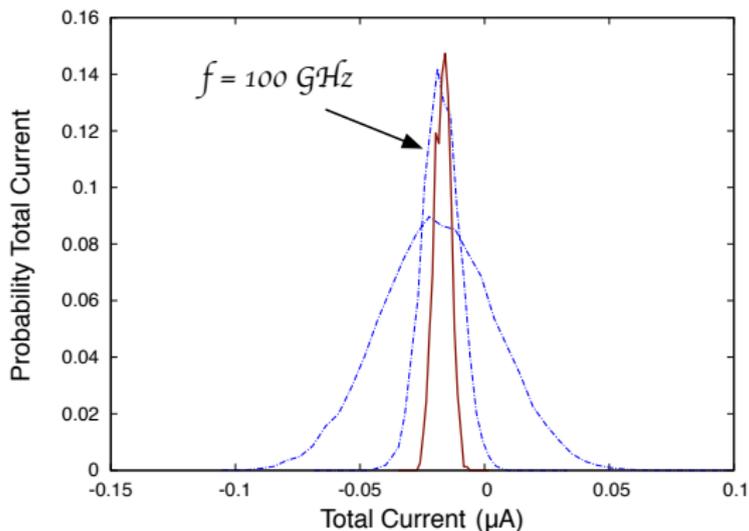
Question: *At which frequency is this additional noise relevant?*



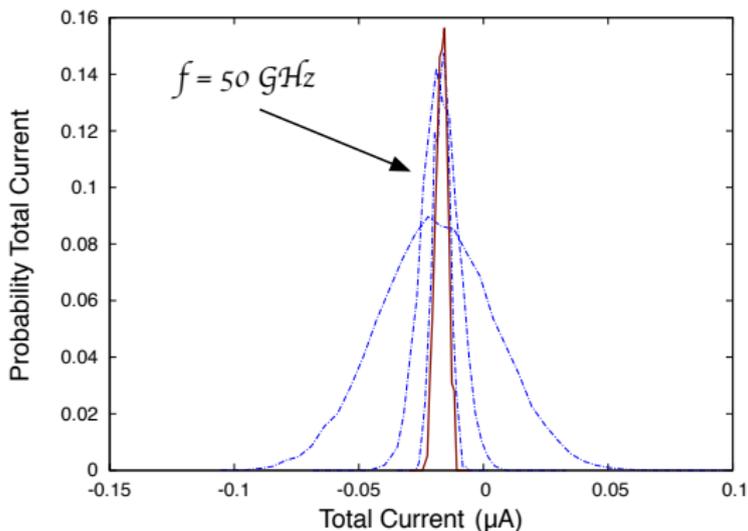
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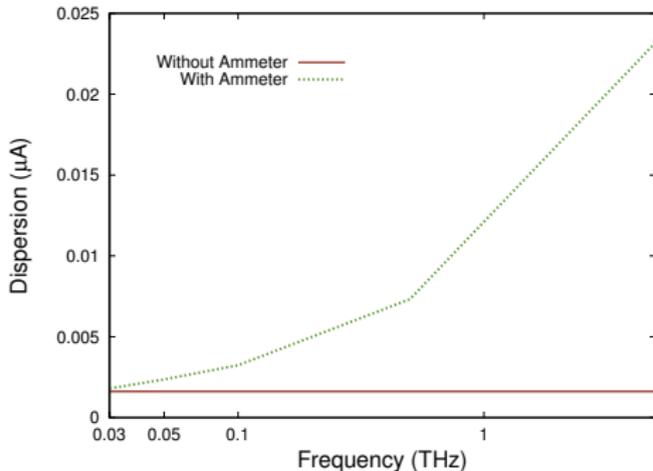
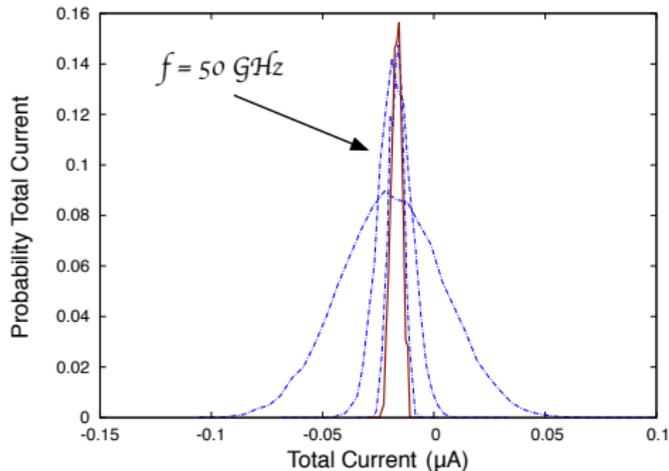
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For the configuration considered $\approx 50 \text{ GHz}$