

# Percolation noise at the metal–insulator transition of nanostructured VO<sub>2</sub> films

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Measurements: (2008 unsuccessful) successful: 2012 – 2014

Zareh



Laszlo

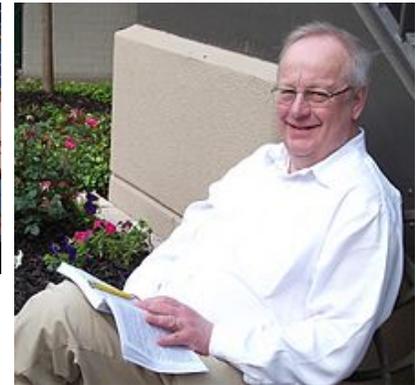
Shuyi



Gunnar



Claes



## Resistance noise at the metal–insulator transition in thermochromic VO<sub>2</sub> films

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(Received 7 November 2014; accepted 30 December 2014; published online 12 January 2015)

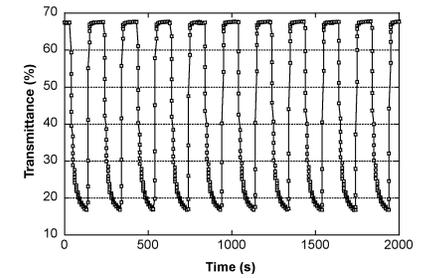
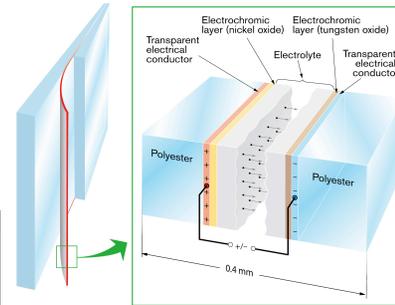
Thermochromic VO<sub>2</sub> films were prepared by reactive DC magnetron sputtering onto heated sapphire substrates and were used to make 100-nm-thick samples that were 10 μm wide and 100 μm long. The resistance of these samples changed by a factor ~2000 in the 50 < T<sub>s</sub> < 70 °C range of temperature T<sub>s</sub> around the “critical” temperature T<sub>c</sub> between a low-temperature semiconducting phase and a high-temperature metallic-like phase of VO<sub>2</sub>. Power density spectra S(f) were extracted for resistance noise around T<sub>c</sub> and demonstrated unambiguous 1/f behavior. Data on S(10 Hz)/R<sub>s</sub><sup>2</sup> scaled as R<sub>s</sub><sup>x</sup>, where R<sub>s</sub> is sample resistance; the noise exponent x was –2.6 for T<sub>s</sub> < T<sub>c</sub> and +2.6 for T<sub>s</sub> > T<sub>c</sub>. These exponents can be reconciled with the Pennetta–Trefán–Reggiani theory [C. Pennetta, G. Trefán, and L. Reggiani, Phys. Rev. Lett. 85, 5238 (2000)] for lattice percolation with switching disorder ensuing from random defect generation and healing in steady state. Our work hence highlights the dynamic features of the percolating semiconducting and metallic-like regions around T<sub>c</sub> in thermochromic VO<sub>2</sub> films.

# VO<sub>2</sub>: applications thermochromic glazing

Energy-Saving Applications. Smart-windows:  
Voltage or temperature controlled transparency

**Transparent**

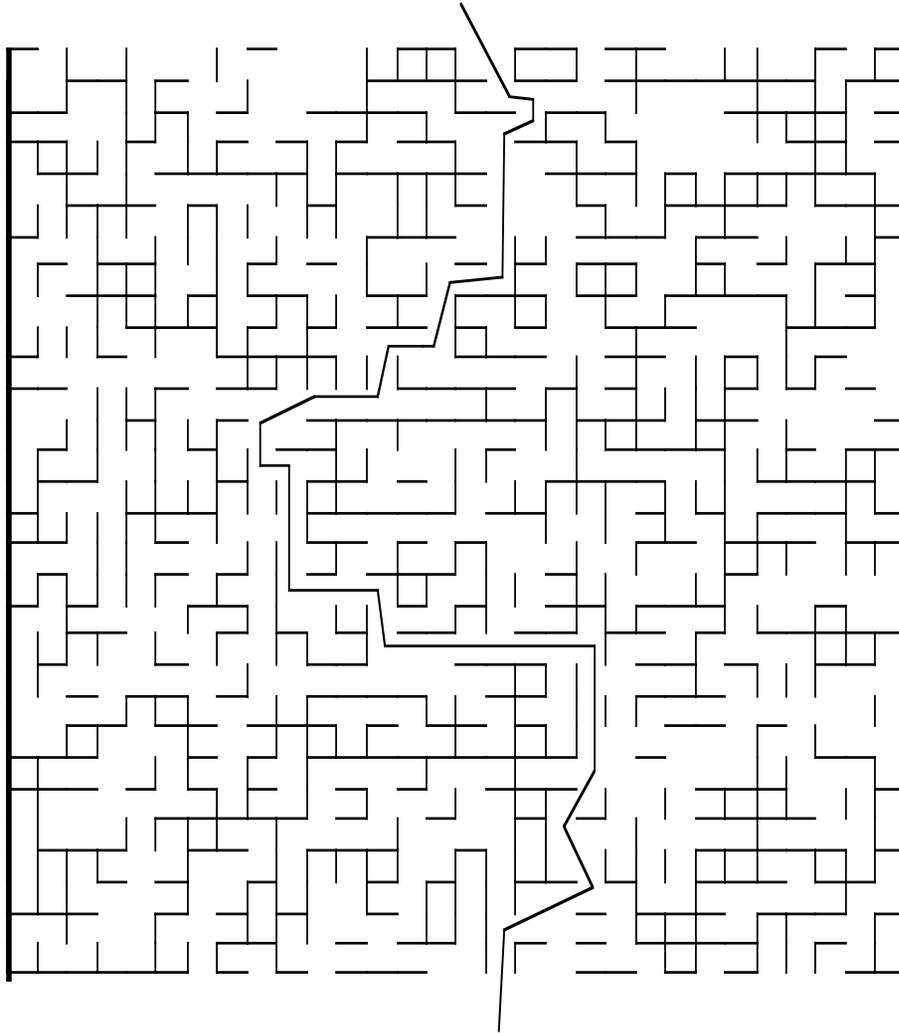
**Darkened**



## VO<sub>2</sub>:

- Thermochemical [F. J. Morin, Phys. Rev. Lett. 3, 34 (1959)];
- Single crystal: first-order metal–insulator transition (MIT) at  $T_c \approx 68$  °C; switching between a low-temperature (monoclinic) *semiconducting* state and a high-temperature (rutile) *metallic-like* state.
- Thin films: the MIT is gradual with *metallic-like* regions growing in extent as the sample temperature  $T_s$  approaches  $T_c$  from below and with *semiconducting* regions disappearing as  $T_s$  becomes increasingly larger than  $T_c$ .

## Lattice percolation



$$R(p) \propto (p - p_c)^{-t}$$

$p$ : filling factor (0 – 1)

$p_c$ : percolation threshold

$t$ : resistivity exponent (1 - 2)

depends on dimension only (in regular lattices)

**because percolation length is scaling in a similar fashion**

$$\frac{S_R(f)}{R^2} \propto R^x$$

resistivity and noise exponent depends on dimension only (in regular lattices) and their absolute value is the same in 2D for conductor-insulator/superconductor transition (due to duality in 2D)

figure from: Z. Gingl, et al, Semicond. Sci. & Technol. 11 (1996) 1770.

Lattice percolation is confirmed in VO<sub>2</sub> films by J. Rozen, et al, Appl. Phys. Lett. 88, 081902 (2006).

$p$ : determined from transparance measurements

$$R(p) \propto (p - p_c)^{-t}$$

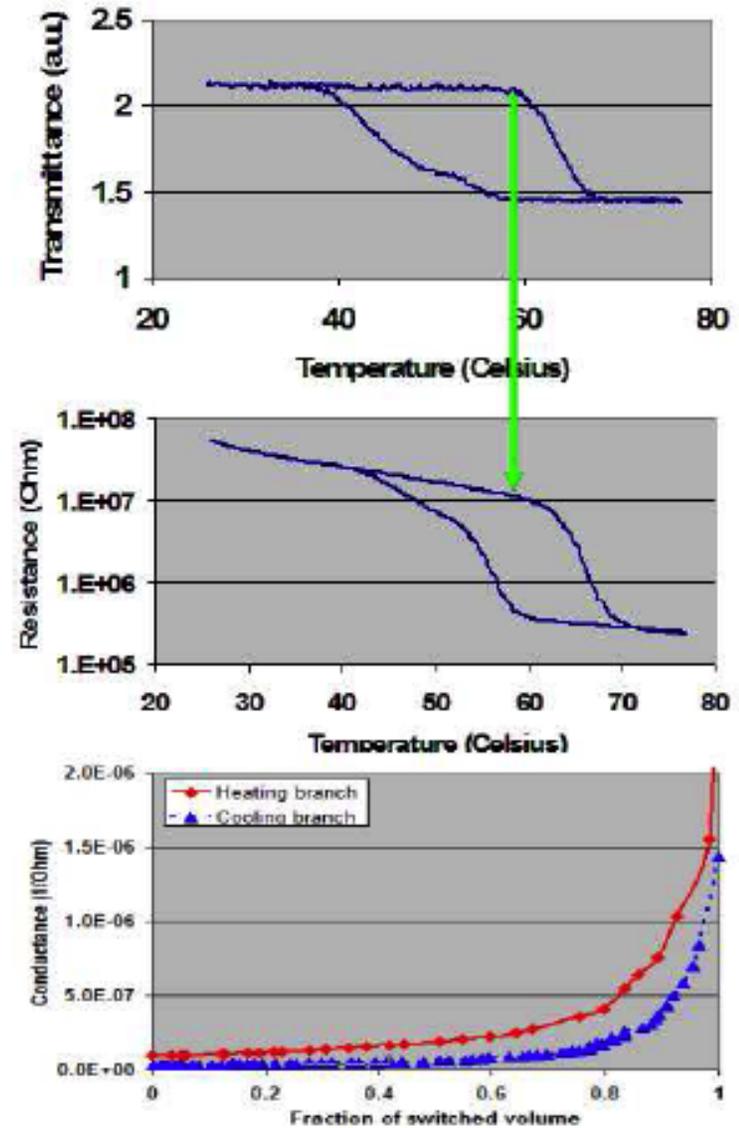
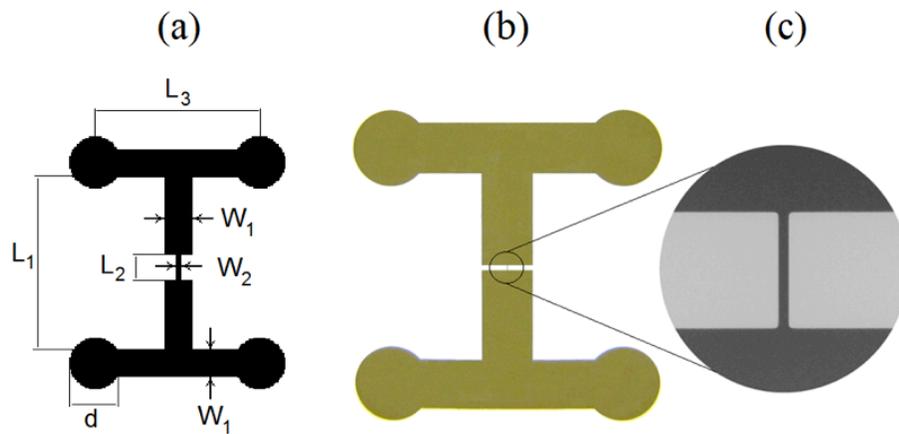


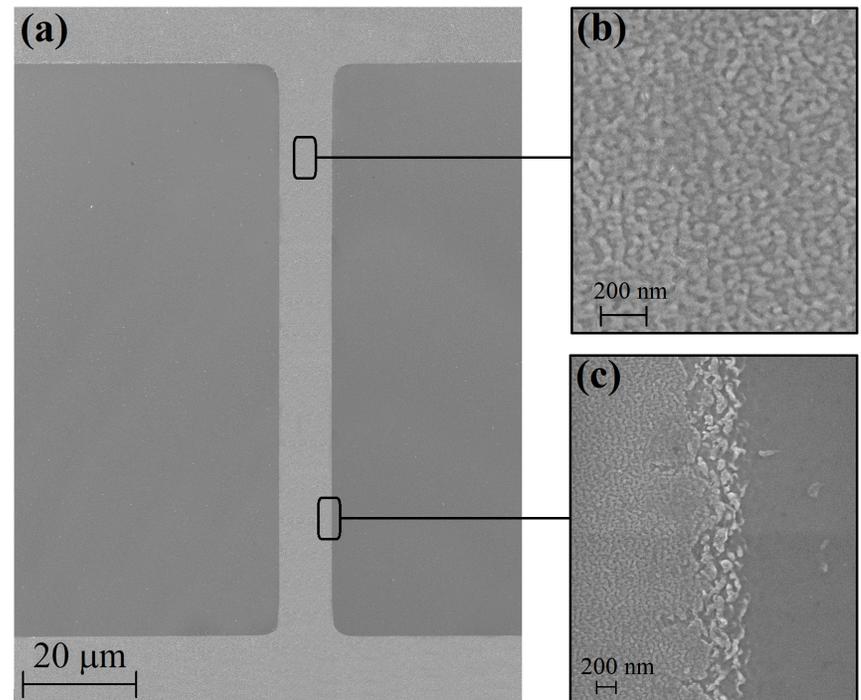
FIG. 3 Transmittance (a) and resistance (b) curves measured simultaneously as a function of temperature. (c) Dependence of the conductance on the volume fraction of the metallic phase deduced from the experimental data.

# Noise experiments.

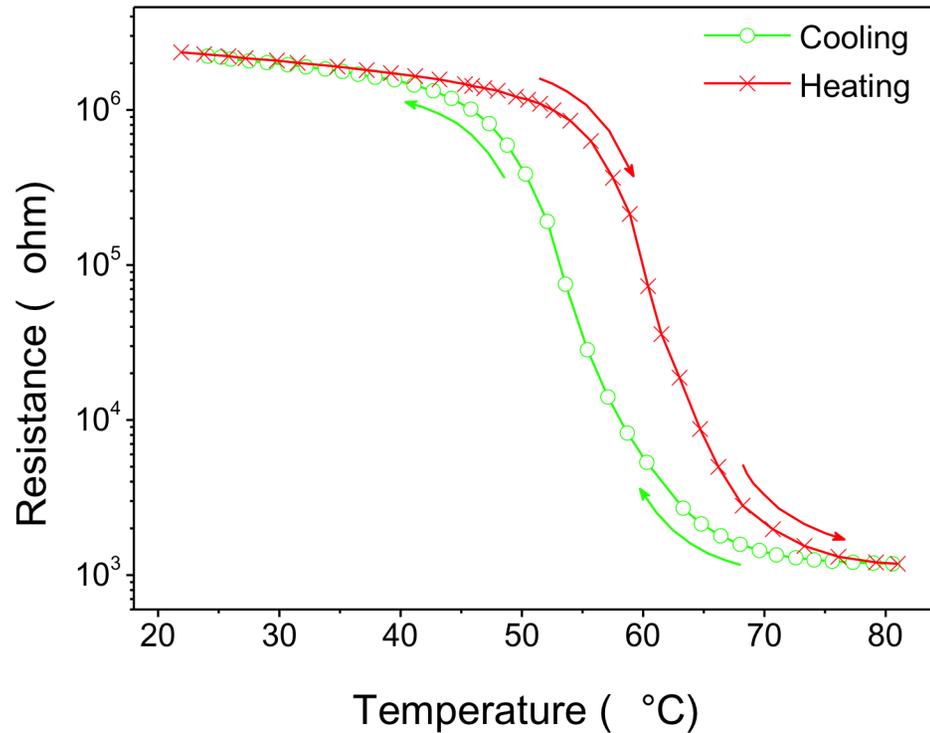
Schematic illustration of a **100-nm**-thick VO<sub>2</sub> sample with a micro-bridge in the middle and with contact pads for four-point electrical measurements. The dimensions are  $L_1 = 5$  mm,  $L_2 = \mathbf{100\ \mu\text{m}}$ ,  $L_3 = 6$  mm,  $d = 1.5$  mm,  $W_1 = 1$  mm and  $W_2 = \mathbf{10\ \mu\text{m}}$ . Panels (b) and (c) are photos of the same structure and of the VO<sub>2</sub> micro-bridge in the encircled region, respectively.



SEM micrographs of the VO<sub>2</sub> micro-bridge



# Resistance measurements



Extraordinary temperature sensitivity. For successful conductance noise measurements, it requires *ultra-low noise temperature control*, which is not available on the market.

We had similar problems in the 1990's with measuring high- $T_c$  superconductor noise.

## New Noise Exponents in Random Conductor-Superconductor and Conductor-Insulator Mixtures

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(Received 29 April 1993)

Time dependent fluctuations of the fraction of normal-conducting part in random resistor-superconductor (RS) and resistor-insulator (RI) networks lead to a novel effect close to the percolation threshold. The normalized noise scales as a function of the resistance with a characteristic exponent  $\lambda$ . The value of  $\lambda$  is different from the value found in classical percolation models but can be related to the resistivity exponent  $s(t)$  of the RS (RI) transition by a simple scaling relation:  $\lambda = 2/s$  ( $2/t$ ). Results of recent experiments on high- $T_c$  superconducting thin films are interpreted in terms of this new effect and a crossover from three to two dimensional percolation behavior is found.

September 19, 1989



Peter

Laszlo

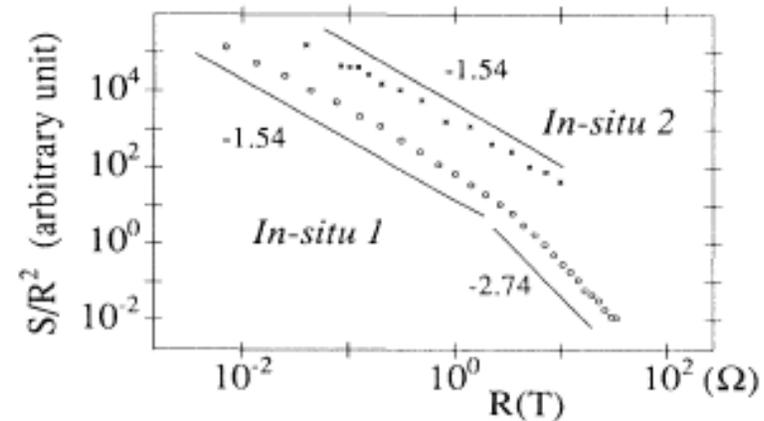


FIG. 1. Normalized noise versus resistance in high-quality high- $T_c$  superconducting thin films in the percolation region. *In situ 1*: sample fabricated by coevaporation. *In situ 2*: sample fabricated by dc magnetron sputtering. The solid lines correspond to slopes predicted by scaling theory.

Ultra-low noise temperature control (originally developed by Per Nordblad, which we modified for the new needs)



- **P. Nordblad**, “Magnetic Anisotropy and Magnetic Phase Transitions of Iron and Manganese Compounds”, in Abstracts of Uppsala Dissertations from the Faculty of Science 556, Acta Universitatis Upsaliensis (Uppsala, Sweden, 1980).

- P. Svedlindh, K. Gunnarsson, P. Nordblad, L. Lundgren, H. Aruga, and A. Ito, Phys. Rev. B 40, 7162 (1989).

- ***Copper thermometer, (DC-) heater and the film sample on the same copper block;***

- ***in vacuum, thus passive thermal relaxation time > 1000 seconds;***

- ***thermometer with 4-point driving/probing arrangement with a differential transformer;***

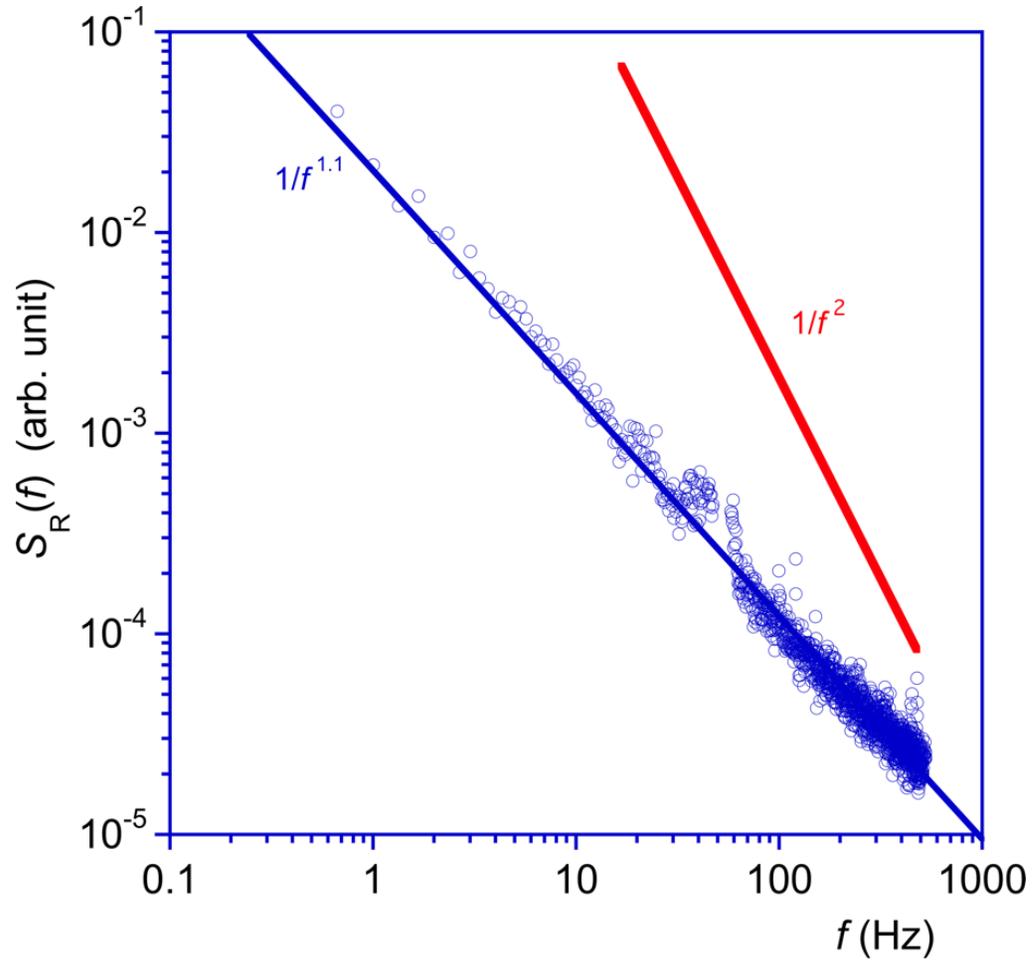
- ***temperature measurement bridge is of high-stability resistors in oil bath***

- ***which is driven by AC, 473Hz to reduce 50 Hz harmonics, generated, filtered/measured by a lockin amplifier***

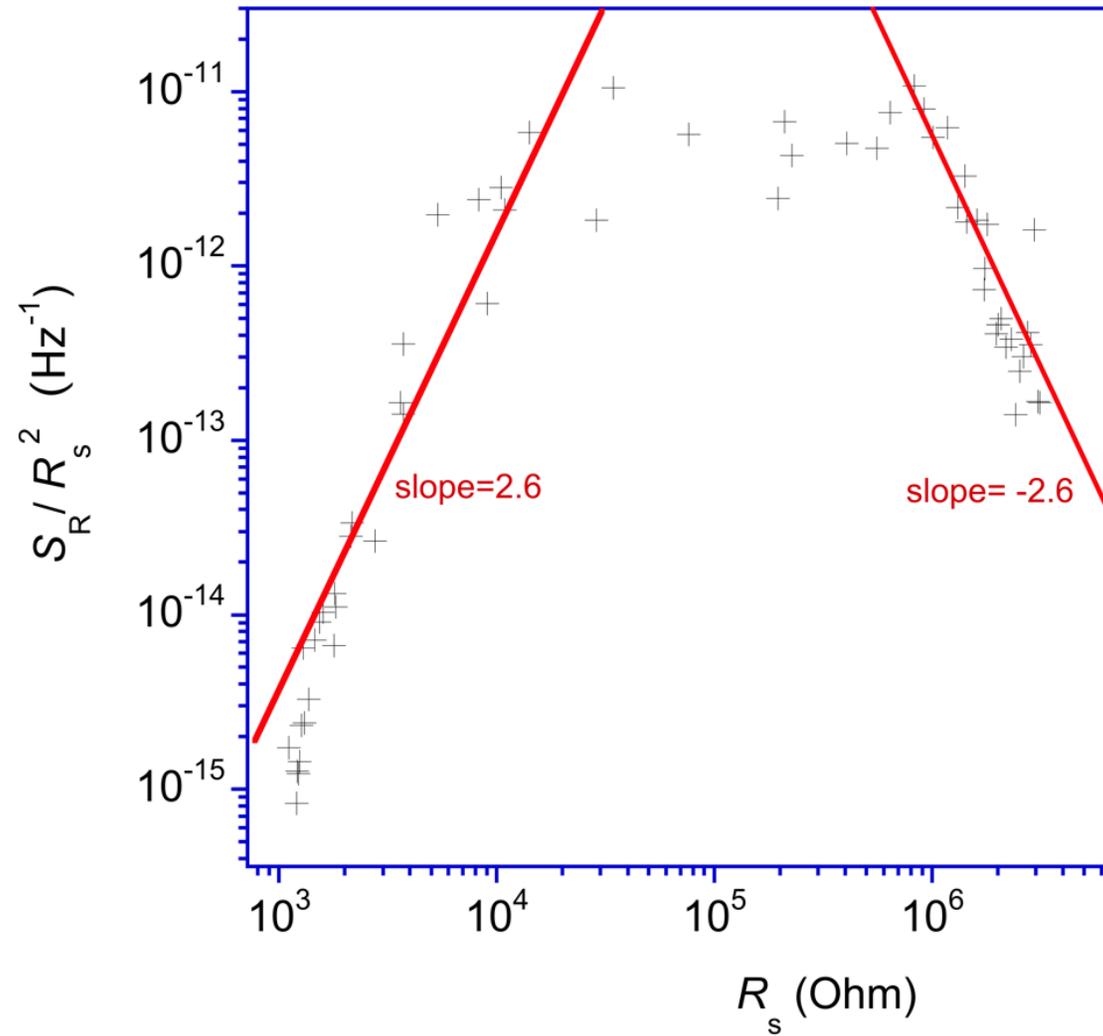
- ***lockin DC output is driving a PID controller, which drives the analog DC heater amplifier.***

**Temperature noise less than  $10^{-9}$  K/Hz<sup>0.5</sup> can be achieved.**

# Conductance noise spectrum and checking for temperature fluctuations ( $1/f^2$ )



# Scaling plot of the normalized noise versus the resistance



Our measured/fitted noise exponents in various high- $T_c$  superconductor films (1989-1994)

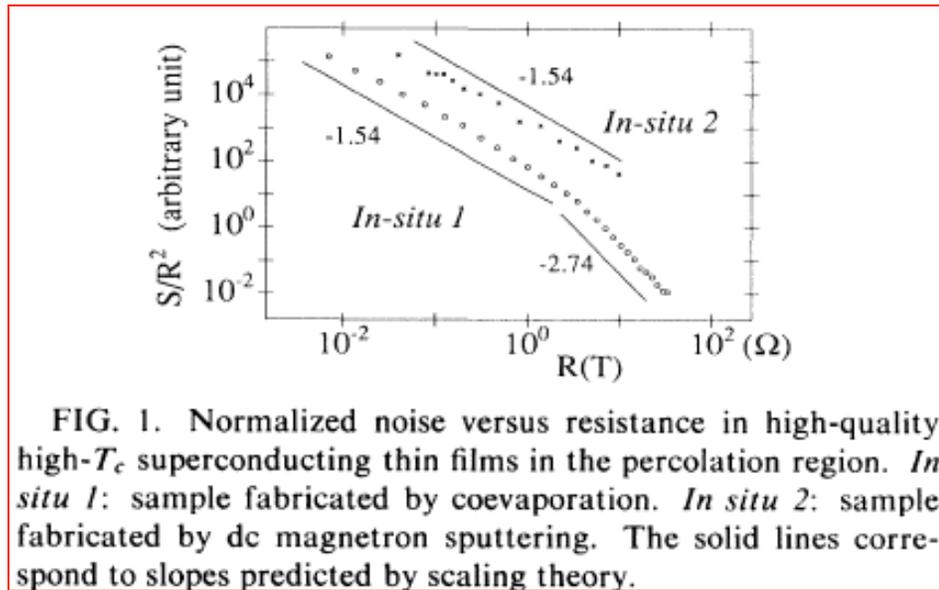
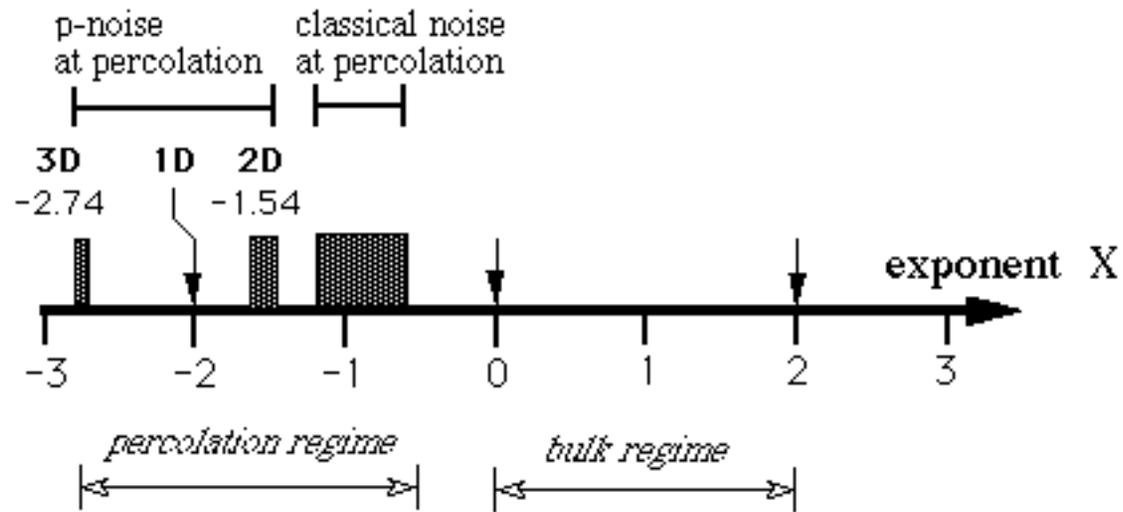
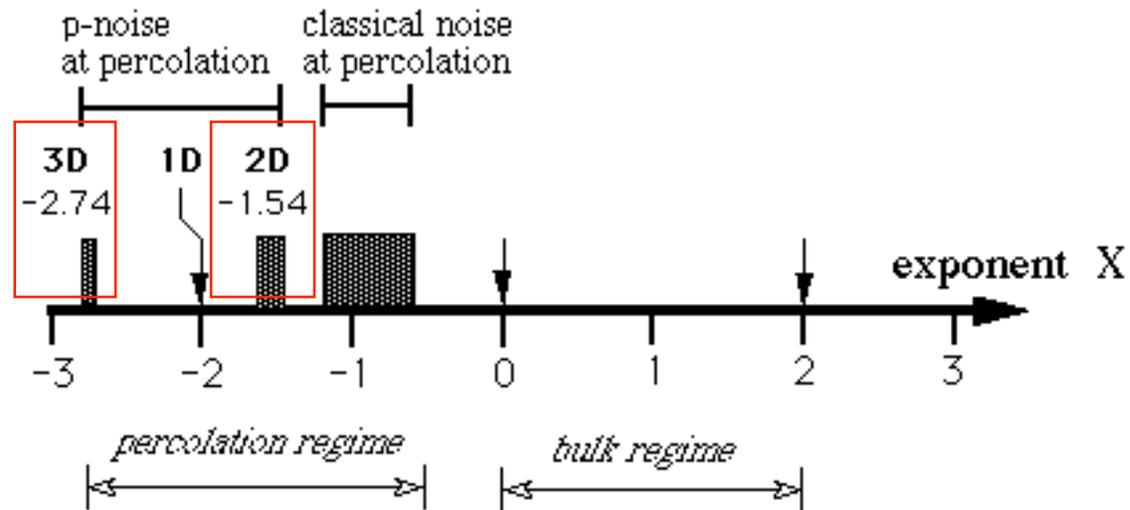


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## Our measured/fitted noise exponents in various high- $T_c$ superconductor films (1989-1994)



The -2.7 exponent would be fine at the high-temperature end however that is the 3D case where duality [P.M. Hui and D. Stroud, Phys. Rev. B 34, 8101 (1986)] does not force the same absolute value in at the low-temperature end as at the high one.

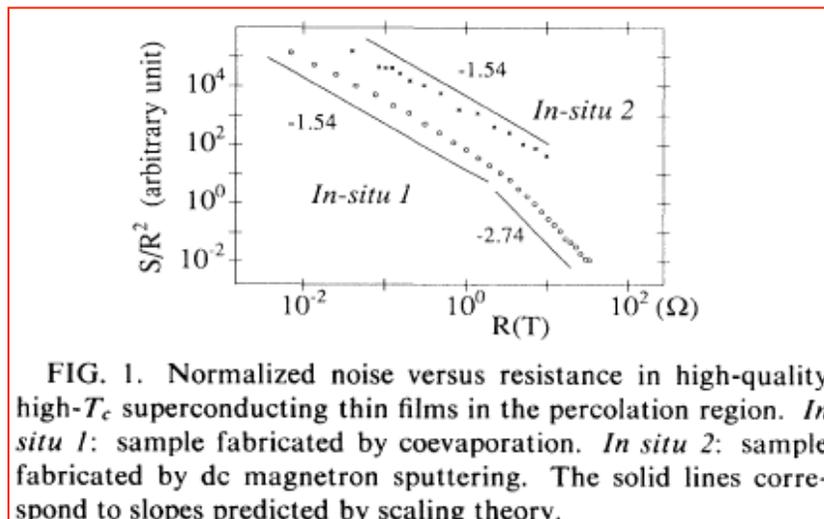


FIG. 1. Normalized noise versus resistance in high-quality high- $T_c$  superconducting thin films in the percolation region. *In situ 1*: sample fabricated by coevaporation. *In situ 2*: sample fabricated by dc magnetron sputtering. The solid lines correspond to slopes predicted by scaling theory.

Kiss-Svedlindh, PRL 1993.

TABLE I. Scaling exponents of the resistance and the normalized noise in random resistor networks. The resistance of the normal-conducting elements fluctuates independently; i.e., the number of noise sources is given by the number of resistors.

|    | RS composite ( $p_r > p_c$ )    |   | RI composite ( $p_r > p_c$ )       |   |
|----|---------------------------------|---|------------------------------------|---|
|    | $R_{rs} \propto (p_r - p_c)^s;$ | $\frac{S_{rs}(f)}{R_{rs}^2} \propto R_{rs}^{-l_{rs}}$ | $R_{ri} \propto (p_r - p_c)^{-t};$ | $\frac{S_{ri}(f)}{R_{ri}^2} \propto R_{ri}^{-l_{ri}}$ |
|    | $s$                             | $l_{rs}$  | $t$                                | $l_{ri}$  |
| 1D | 1                               | 1   | ...                                | ...   |
| 2D | $1.297 \pm 0.07$                | $0.86 \pm 0.02$                                       | $1.297 \pm 0.07$                   | $0.86 \pm 0.02$                                       |
| 3D | $0.73 \pm 0.011$                | $0.9 \pm 0.32$  | $1.96 \pm 0.1$                     | $0.80 \pm 0.1$  |

*Kiss-Svedlindh, p-fluctuations, percolation noise model, 1993*

TABLE II. Comparison of classical ( $l$ ) and new ( $\lambda$ ) noise exponents.

|    | RS composite ( $p_r > p_c$ )                      |   | RI composite ( $p_r > p_c$ )                           |   |
|----|---|---|--|---|
|    | $\frac{S_{rs}(f)}{R_{rs}^2} \propto R^{-l_{rs}};$ | $\frac{S_{rs}(f)}{R_{rs}^2} \propto R_{rs}^{-\lambda_{rs}}$ | $\frac{S_{ri}(f)}{R_{ri}^2} \propto R_{ri}^{-l_{ri}};$ | $\frac{S_{ri}(f)}{R_{ri}^2} \propto R_{ri}^{-\lambda_{ri}}$ |
|    | $l_{rs}$  | $\lambda_{rs} = 2/s$  | $l_{ri}$   | $\lambda_{ri} = 2/t$  |
| 1D | 1   | 2   | ...  | ...   |
| 2D | $0.86 \pm 0.02$                                   | $1.54 \pm 0.09$   | $0.86 \pm 0.02$  | $1.54 \pm 0.09$   |
| 3D | $0.9 \pm 0.32$                                    | $2.74 \pm 0.04$   | $0.80 \pm 0.1$   | $1.02 \pm 0.05$   |

## Scaling Law of Resistance Fluctuations in Stationary Random Resistor Networks

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(Received 24 July 2000)

In a random resistor network we consider the simultaneous evolution of two competing random processes consisting in breaking and recovering the elementary resistors with probabilities  $W_D$  and  $W_R$ . The condition  $W_R > W_D/(1 + W_D)$  leads to a stationary state, while in the opposite case, the broken resistor fraction reaches the percolation threshold  $p_c$ . We study the resistance noise of this system under stationary conditions by Monte Carlo simulations. The variance of resistance fluctuations  $\langle \delta R^2 \rangle$  is found to follow a scaling law  $|p - p_c|^{-\kappa_0}$  with  $\kappa_0 = 5.5$ . The proposed model relates quantitatively the defectiveness of a disordered media with its electrical and excess-noise characteristics.



The **Pennetta-Trefan-Reggiani** model of "dynamical percolation" with microscopic damage and healing processes with separate rates in 2D produces 2.6 exponents in the steady-state at low-temperature.

In 2D, *due to duality* [P. M. Hui and D. Stroud, Phys. Rev. B 34, 8101 (1986)], the same absolute exponent value holds in the high-temperature limit with negative sign.

tain, in the disordered network regime,

$$\frac{\langle \delta R^2 \rangle}{\langle R \rangle^2} \sim \langle R \rangle^{-s} \quad (8)$$

with  $s = 2.6$  as reported in Fig. 6.

in the percolation (scaling) region:

$$\frac{S_R(f)}{R^2} \propto R^x \quad \text{with} \quad x \approx \pm 2.6$$

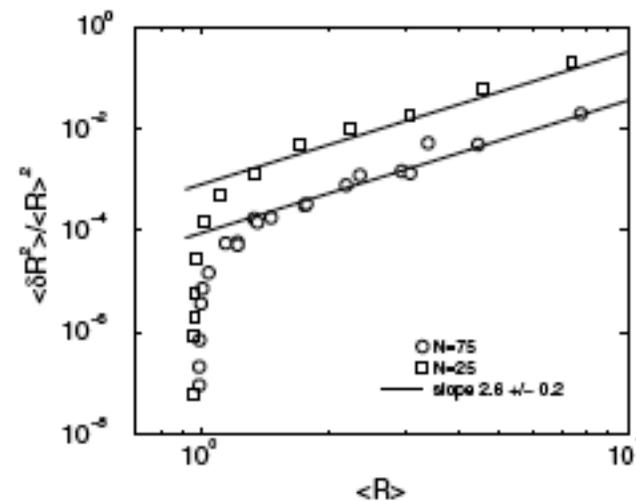
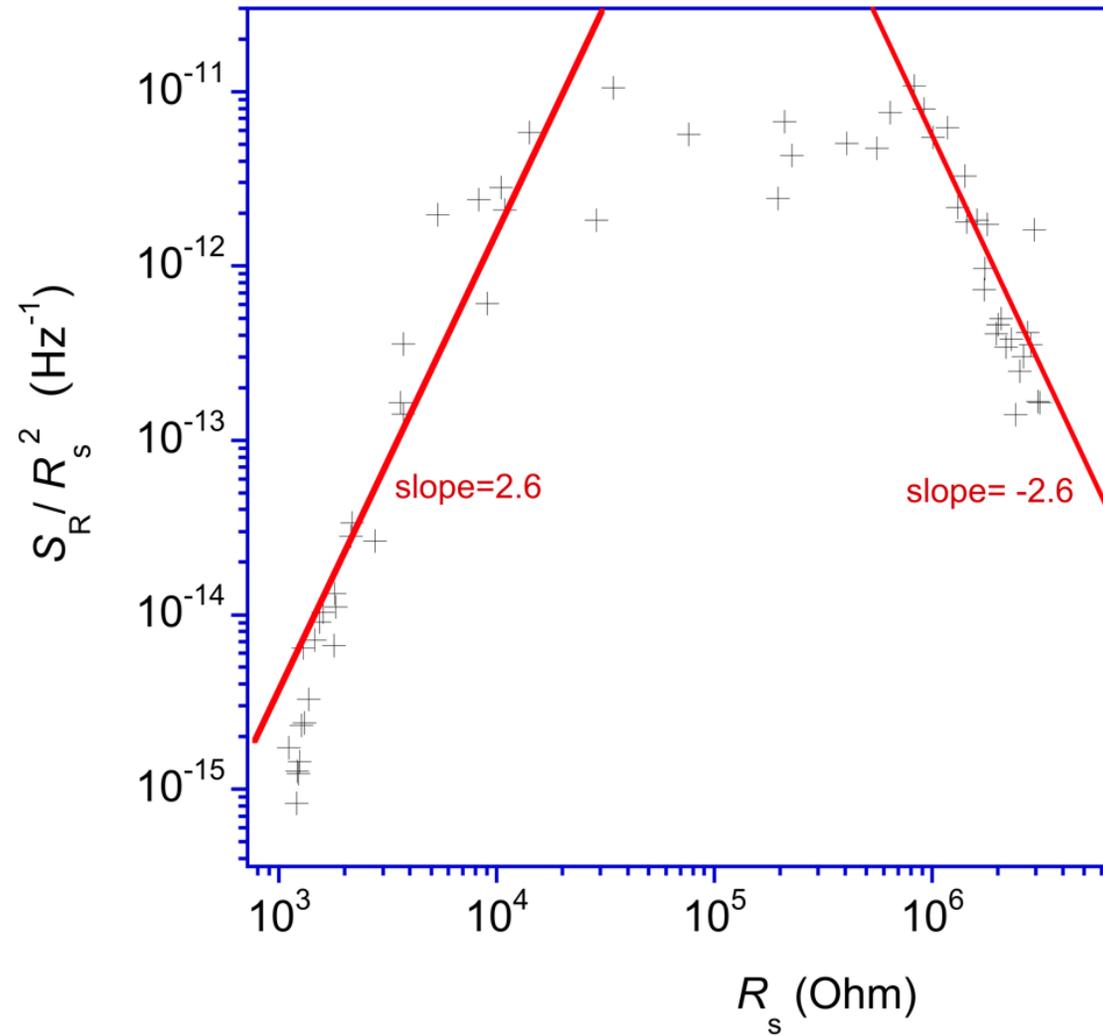


FIG. 6. Resistance noise normalized to the square of the average RRN resistance  $\langle \delta R^2 \rangle / \langle R \rangle^2$  as a function of the average resistance  $\langle R \rangle$ .

# Scaling plot of the normalized noise versus the resistance



## UPoN!!!

- In the PTR model, the noise is not inherent in the resistance but comes from the switching
- PTR see Lorentzian spectra. How do we get  $1/f$  noise?
- Hierarchy of switching time constants?
- Why is the spectrum  $1/f$  in the whole temperature range?
- Perhaps another model is relevant?