

Brownian motors in the micro-scale domain: Enhancement of efficiency by noise

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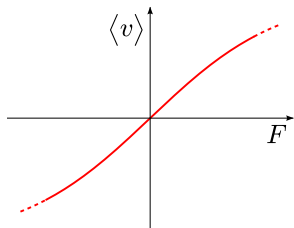
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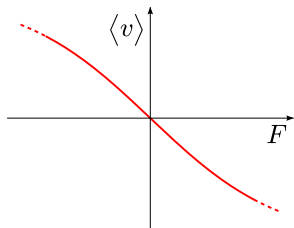
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Le Chatèlier-Braun principle

$$\langle v \rangle = \mu F.$$



Typical response



Absolute negative mobility

Minimal model exhibiting ANM expressed in the dimensionless variables

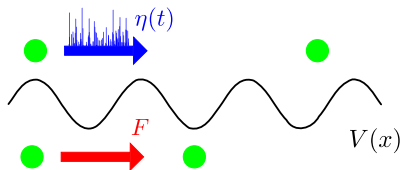
$$\ddot{x} + \gamma \dot{x} = -V'(x) + a \cos(\omega t) + F + \sqrt{2\gamma D_G} \xi(t).$$

We replace the constant static load F with the random force $\eta(t)$

$$F \rightarrow \eta(t), \quad \text{assuming} \quad \langle \eta(t) \rangle = F.$$

Fundamental question

Can noise $\eta(t)$ induce more effective transport than the constant force F ?



Nonequilibrium noise

Generalized white Poissonian noise

$$\eta(t) = \sum_{i=1}^{n(t)} z_i \delta(t - t_i),$$

where $n(t)$ is Poissonian counting process

$$\text{Pr}\{n(t) = k\} = \frac{(\lambda t)^k}{k!} e^{-\lambda t}.$$

The process $\eta(t)$ presents white noise of finite mean and a covariance given by

$$\langle \eta(t) \rangle = \lambda \langle z_i \rangle, \quad \langle \eta(t) \eta(s) \rangle - \langle \eta(t) \rangle \langle \eta(s) \rangle = 2D_P \delta(t - s).$$

Its intensity reads

$$D_P = \frac{\lambda \langle z_i^2 \rangle}{2}.$$

Distribution of the amplitudes $\{z_i\}$ of the δ -kicks

$$\rho(z) = \zeta^{-1} \theta(z) e^{-z/\zeta}, \quad \langle z_i^k \rangle = k! \zeta^k, \quad k = 1, 2, \dots$$

The mean value and the intensity of white Poissonian shot noise

$$\langle \eta(t) \rangle = \lambda \zeta = \sqrt{D_P \lambda} \geq 0, \quad D_P = \lambda \zeta^2.$$

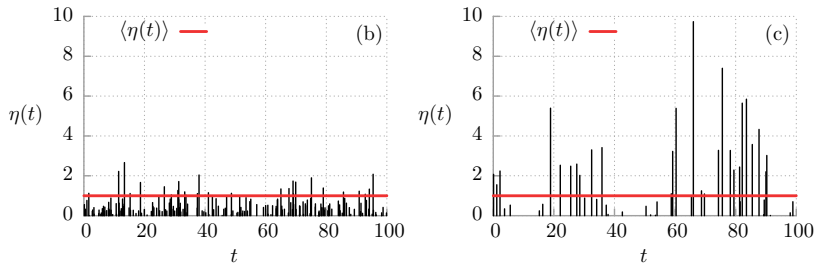


Figure : In (b): $\lambda = 2$, $D_P = 0.5$; in (c): $\lambda = 0.5$, $D_P = 2$.

- Average velocity

$$\lim_{t \rightarrow \infty} \langle \dot{x}(t) \rangle = \langle v \rangle + v_{\omega}(t) + v_{2\omega}(t) + \dots$$

$$\langle v \rangle = \lim_{t \rightarrow \infty} \frac{\omega}{2\pi} \int_t^{t+2\pi/\omega} ds \langle v(s) \rangle .$$

- Velocity fluctuations

$$\sigma_v^2 = \langle v^2 \rangle - \langle v \rangle^2, \quad v(t) \in [\langle v \rangle - \sigma_v, \langle v \rangle + \sigma_v].$$

- Stokes efficiency, $P_{out} = \gamma \langle v \rangle^2$, $P_{in} = \gamma [\langle v \rangle^2 + \sigma_v^2 - D_G]$

$$\varepsilon_S = \frac{P_{out}}{P_{in}} = \frac{\langle v \rangle^2}{\langle v \rangle^2 + \sigma_v^2 - D_G} .$$

Average velocity

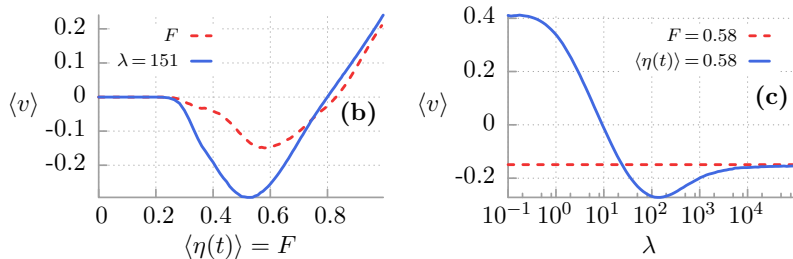


Figure : Anomalous transport regime. Parameters are $a = 8.95$, $\omega = 3.77$, $\gamma = 1.546$, $D_G = 0.001$, $\lambda = 151$.

Velocity fluctuations

$$\sigma_v^2 = \langle v^2 \rangle - \langle v \rangle^2.$$

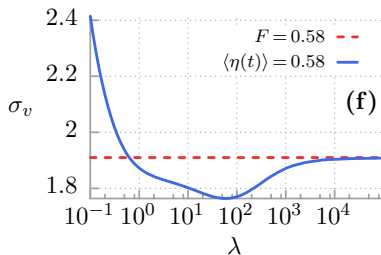
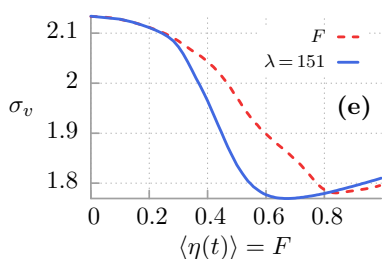


Figure : Anomalous transport regime. Parameters are $a = 8.95$, $\omega = 3.77$, $\gamma = 1.546$, $D_G = 0.001$, $\lambda = 151$.

$$\varepsilon_S = \frac{\langle v \rangle^2}{\langle v \rangle^2 + \sigma_v^2 - D_G}$$

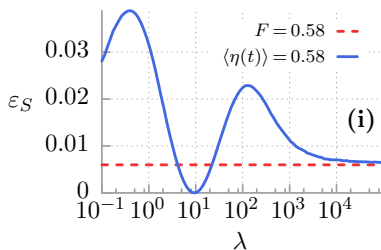
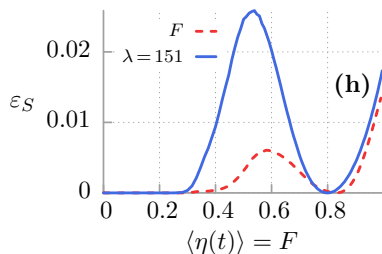


Figure : Anomalous transport regime. Parameters are $a = 8.95$, $\omega = 3.77$, $\gamma = 1.546$, $D_G = 0.001$, $\lambda = 151$.

Fundamental question

Can noise $\eta(t)$ induce more effective transport than the constant force F ?

Answer

Yes, the Brownian motor can move much faster, its velocity fluctuations are much smaller and the motor efficiency increases several times in both normal and absolute negative mobility regimes.

**New operating principle: consider replacing
the constant force by nonequilibrium noise!**

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